

LEPTOGENESIS AND THE GRAVITINO PROBLEM

- Status of Leptogenesis
- Connection with ν -Masses
- The Gravitino Problem
- $\Omega_{DM} \approx \Omega_{3/2}$

I. Status of Leptogenesis

- (simple) thermal leptogenesis:

Fukuyama, Yanagida '86

add ν_{Ri} , $i=1-3$, to SM, seesaw mechanism gives mass eigenstates $\nu_i \approx \nu_{Li} + \nu_{Ri}^c$,

$N_i \approx \nu_{Ri} + \nu_{Ri}^c$, with

$$\underbrace{m_N \approx M, \quad m_\nu \approx -m_D \frac{1}{M} m_D^T,}_{\Delta L = 2}$$

$$\Delta L = 2$$

Dirac neutrino mass matrix $m_D = h \nu$,

$\nu = 174 \text{ GeV}$, $M \gg m_D$, $m_\nu \ll m_D$; CP-violating decays of heavy Majorana neutrinos,

$$\Gamma(N_i \rightarrow LH) \neq \Gamma(N_i \rightarrow L^c H^c),$$

generate lepton asymmetry ΔL , which is partially converted to baryon asymmetry ΔB by sphaleron processes (KRS, $\Delta B = \Delta L = 3$); in therm. equilibrium:

$$T_{EW} \sim 100 \text{ GeV} < T < T_{\text{sph}} \sim 10^{12} \text{ GeV}$$

Thermal leptogenesis is an attractive mechanism for baryogenesis:

- (1) baryon asymmetry determined by neutrino properties
- (2) baryon asymmetry consistent with neutrino masses

Quantitative Analysis

(\rightarrow WB, Di Bari, Plumacher, Ann. Phys. '05)

from nucleosynthesis, with $T_{\text{BBN}} \sim 1 \text{ MeV}$,

$$M_{\text{B}}^{\text{BBN}} = \frac{\mu_{\text{B}}}{\mu_{\gamma}} = (2.6 - 6.2) \times 10^{-10}$$

from microwave background, with $T_{\text{CMB}} \sim 1 \text{ eV}$,

$$M_{\text{B}}^{\text{CMB}} = (6.1^{+0.3}_{-0.2}) \times 10^{-10}$$

impressive test of standard cosmological model \rightarrow stringent bounds on ν -masses

Theory :

$$\eta_B = 0.01 \epsilon_i \kappa_f \quad (N_i \rightarrow l\phi)$$

CP asymmetry

Davidson,
Barra

$$\epsilon_i^{\max} = \frac{3}{16\tilde{m}} \frac{M_i \mu_{atm}}{v^2} \quad \mu_{atm} = (\Delta m_{atm}^2)^{1/2}$$

$$\approx 10^{-6} \left(\frac{M_i}{10^{10} \text{GeV}} \right) \left(\frac{\mu_{atm}}{0.05 \text{eV}} \right)$$

efficiency factor

$$\kappa_f = (2 \pm 1) \times 10^{-2} \left(\frac{0.01 \text{eV}}{\tilde{m}_i} \right)^{1.1 \pm 0.1}$$

$$m_i \leq \tilde{m}_i \lesssim m_3$$

light neutrino masses :

$$10^{-3} \text{eV} \leq m_i \leq 0.1 \text{eV}$$

cosmology
 $\Sigma m_i \sim 0.10$

initial conditions

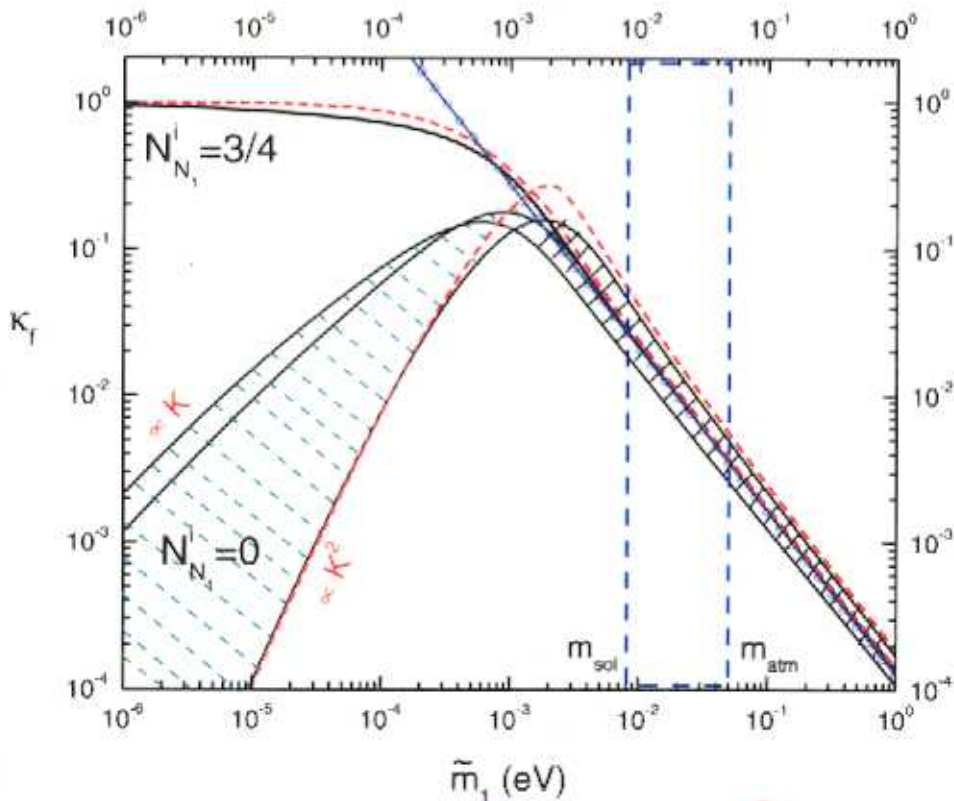
heavy neutrino masses :

$$M_i \gtrsim 10^9 \text{GeV}, \quad T_R \gtrsim 10^9 \text{GeV}$$

INPUT: seesaw (ν_R) ; no triplets, mass degenerates

Efficiency factor

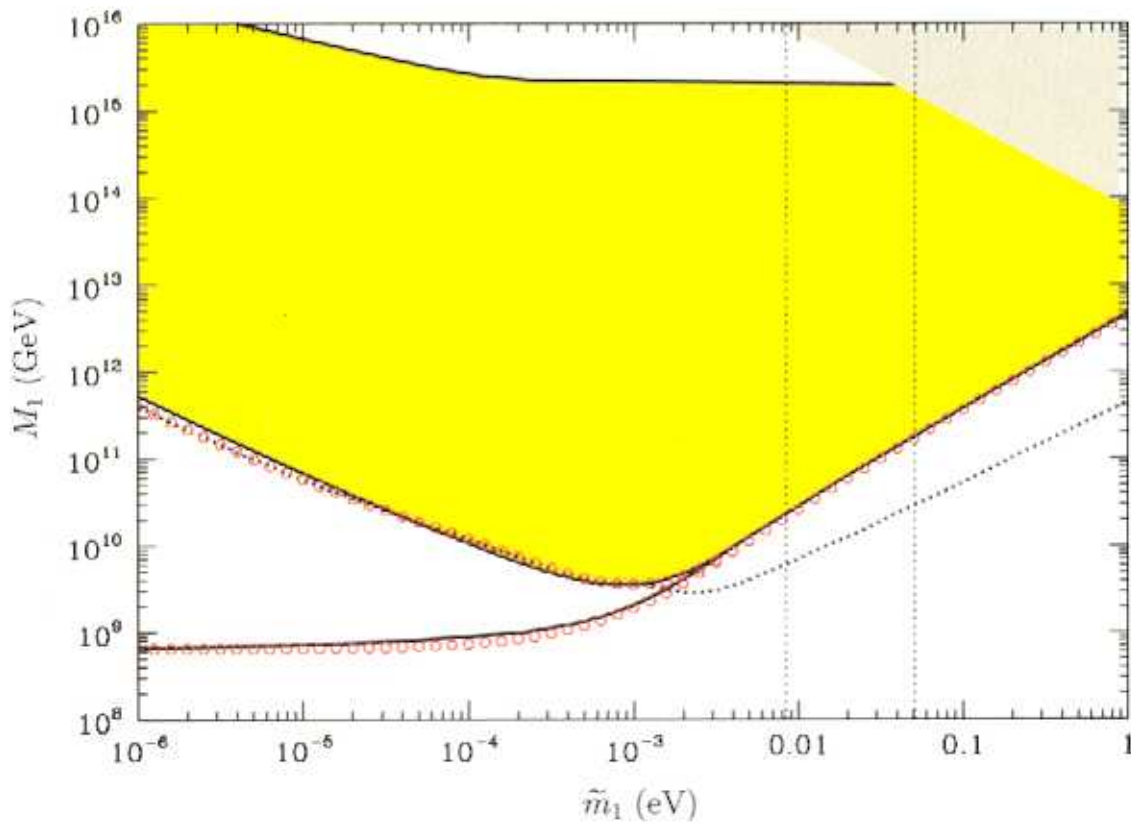
$$\eta_B = d \varepsilon, \kappa_f \approx 10^{-2} \varepsilon, \kappa_f$$



\uparrow \uparrow
 neutrino mass window

$N_{N_1}^{I, \nu} = 0$: $\propto \kappa^2$, only (inverse) decays
 $\propto \kappa$, (inverse) decays + scatterings

Bounds on M_1 and T_B



— M_1^{min}
..... T_B^{min}

II. Connection with ν -masses : 4 examples

(i) SO(10) orbifold GUT in 6D

Aseba, W.B., Coni '03

symmetry breaking :

$$G_{SM} = SU(3) \times SU(2)_L \times U(1)_Y \times U(1)_X \subset SO(10)$$

$$= SU(5) \times U(1)_X \supset SU(4) \times SU(2)_L \times SU(2)_R$$

quarks and leptons are mixtures of
brane and bulk states ; characteristic
predictions for fermion mass matrices :

$$m^u \sim m_N, \quad m^D \sim m^e \sim m^d$$

neutrino sector :

$$M_1 \sim 10^{10} \text{ GeV}, \quad M_2 \sim 3 \times 10^{12} \text{ GeV}, \quad M_3 \sim 10^{15} \text{ GeV}$$

$$m_3 \sim 0.05 \text{ eV}, \quad m_2 \sim m_1 \sim 0.01 \text{ eV}, \quad \theta_{13} \sim 0.1$$

$$\tilde{m}_1 \sim 0.01 \text{ eV}, \quad \epsilon_1 \sim 0.1 \frac{M_1}{M_3} \sim 10^{-6} \quad \text{LG O.K.}$$

factors
O(1)
unknown

$$\text{note : } \frac{m_1}{m_3} \sim \left(\frac{m_d}{m_b} \right)^2 \frac{m_t}{m_u} \sim 0.1$$

(ii) Bi-large neutrino mixing, SO(10) SuSy GUT in 4D

Drumšek, Raby '05

SO(10) broken to G_{SM} by Higgs mechanism;
family symmetry

$$D_3 \times U(1) \times Z_2 \times Z_3$$

strongly restricts fermion mass matrices

$$m^u, m_N, m^D, m^e, m^d$$

χ^2 fit to 20 observables; 14 parameters,
6 predictions; additional prediction:
baryon asymmetry, sign correct!

neutrino sector:

$$M_1 = 1 \times 10^{10} \text{ GeV}, M_2 = -9 \times 10^{11} \text{ GeV}, M_3 = 6 \times 10^{13} \text{ GeV},$$

$$m_3 = 0.05 \text{ eV}, m_2 = 0.01 \text{ eV}, m_1 = 0.004 \text{ eV},$$

$$\theta_{13} = 0.05; \text{ LG: } \epsilon_1 = -1.6 \times 10^{-7},$$

$$\tilde{m}_1 \approx 0.2 \text{ eV} \rightarrow M_B = 1.5 \times 10^{-12} \text{ (?)}$$

(iii) Maximal atmospheric ν -mixing and leptogenesis

Grimus, Lavoue '04

experiment : $\Theta_{23} = 45^\circ$?

What is the underlying symmetry ?

interesting possibility : μ - τ interchange symmetry ; prediction :

$$\Theta_{13} = 0$$

models with family symmetry :

$$(A) \quad Z_2 \quad , \quad (B) \quad D_4$$

compatibility with leptogenesis ?

(B) : disfavoured

(A) : predictions for M_1 , m_1

$$10^{11} \text{ GeV} \lesssim M_1 \lesssim 10^{12} \text{ GeV}$$

$$10^{-3} \text{ eV} \lesssim m_1 \lesssim 10^{-2} \text{ eV}$$

$\Rightarrow m_1 \gtrsim 10^{-3} \text{ eV}$, independence of initial conditions

(iv) Nonthermal leptogenesis

Alm, Kolb '05

recent example: 'Instant leptogenesis',
 in reheating after inflation; inflaton (ϕ)
 Higgs (H) coupling ($g = O(1)$):

$$L = -\frac{1}{2} g^2 \phi^2 H^\dagger H,$$

$m_H = g|\phi|$ oscillates, $m_H > m_N$, $m_H < m_N$,
 $\Delta L (\rightarrow \Delta B)$ from CP violating decays

$$N \rightarrow Hl (H\bar{l}), \quad H \rightarrow Nl (N\bar{l})$$

constraints on neutrino masses:

$$10^{11} \text{ GeV} (10^{12} \text{ GeV}) < M_1 < 10^{13} \text{ GeV}$$

$$3 \times 10^{-4} \text{ eV} < \hat{m}_1 < 1 \text{ eV} (0.1 \text{ eV})$$

reheating temperature: $T_R > 10^{10} \text{ GeV}$

(v) Resonant (TeV scale) leptogenesis (?)

degeneracy: $\frac{M_2 - M_1}{M_1} \sim 10^{-10}$

Pilaftsis,
 Giudice

CP asymmetry (\rightarrow hep-ph/0511248,
 Anisimov, Branco, Pliumachov)

III. The Gravitino Problem

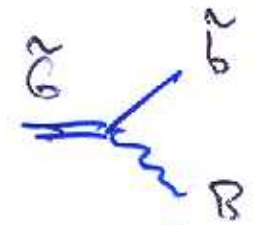
gravitino couplings are universal, $\propto \frac{1}{M_P}$; the mass scale of SUSY breaking, i.e. the gravitino mass is presently unknown:

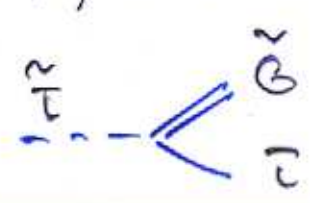
anomaly mediation	$\sim 10 \text{ TeV} \dots 100 \text{ TeV}$
gravity	$\sim 100 \text{ GeV} \dots 1 \text{ TeV}$
gaugino	$\sim 10 \text{ GeV} \dots 100 \text{ GeV}$
gauge	$\sim 10 \text{ eV} \dots 10 \text{ GeV}$

(split SUSY ...)

\rightarrow $10 \text{ eV} \lesssim m_{3/2} \lesssim 100 \text{ TeV} \quad (< M_P)$

quasi-stable particles:

$\tau(\tilde{G} \rightarrow \tilde{b} B) \approx 4 \cdot 10^8 \left(\frac{100 \text{ GeV}}{m_{\tilde{G}}}\right)^3 \text{ s}$


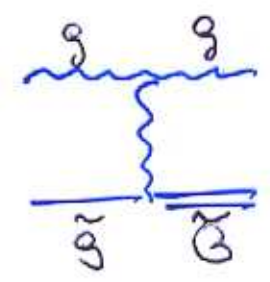
$\tau(\tilde{\tau} \rightarrow \tilde{G} \tau) \approx 8 \cdot 10^3 \left(\frac{m_{3/2}}{1 \text{ GeV}}\right)^2 \left(\frac{150 \text{ GeV}}{m_{\tilde{\tau}}}\right)^5 \text{ s}$


Cosmological 'problems'

severe constraints from BBN on quasi-stable heavy particles, overclosure of universe, $\Omega_{DM} > 1$ (Weinberg '82; Khlopov, Linde; Ellis, Gium, Nanopoulos '84; ...)

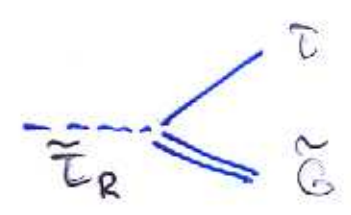
thermal production:

$$\Omega_{3/2}^{TH} h^2 \propto \frac{1}{M_{pl}^2} \left(1 + \frac{m_{\tilde{g}}^2}{3m_{3/2}^2}\right) T_R$$



non-thermal production (\tilde{G} stable):

$$\Omega_{3/2}^{NT} h^2 \propto m_{3/2} m_{\tilde{L}_R} \approx \frac{m_{\tilde{L}_R}}{LSP}$$

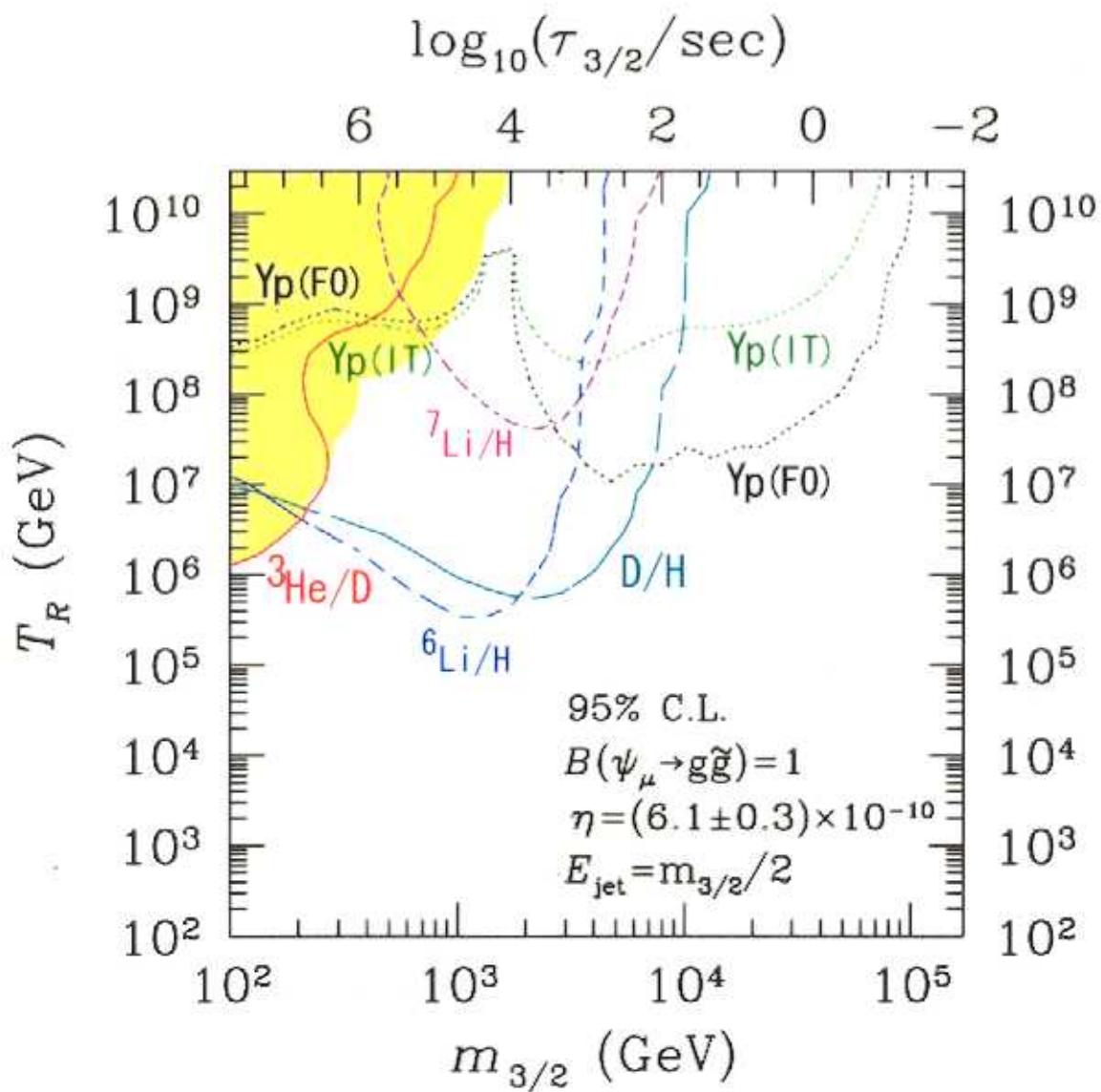


constraints:

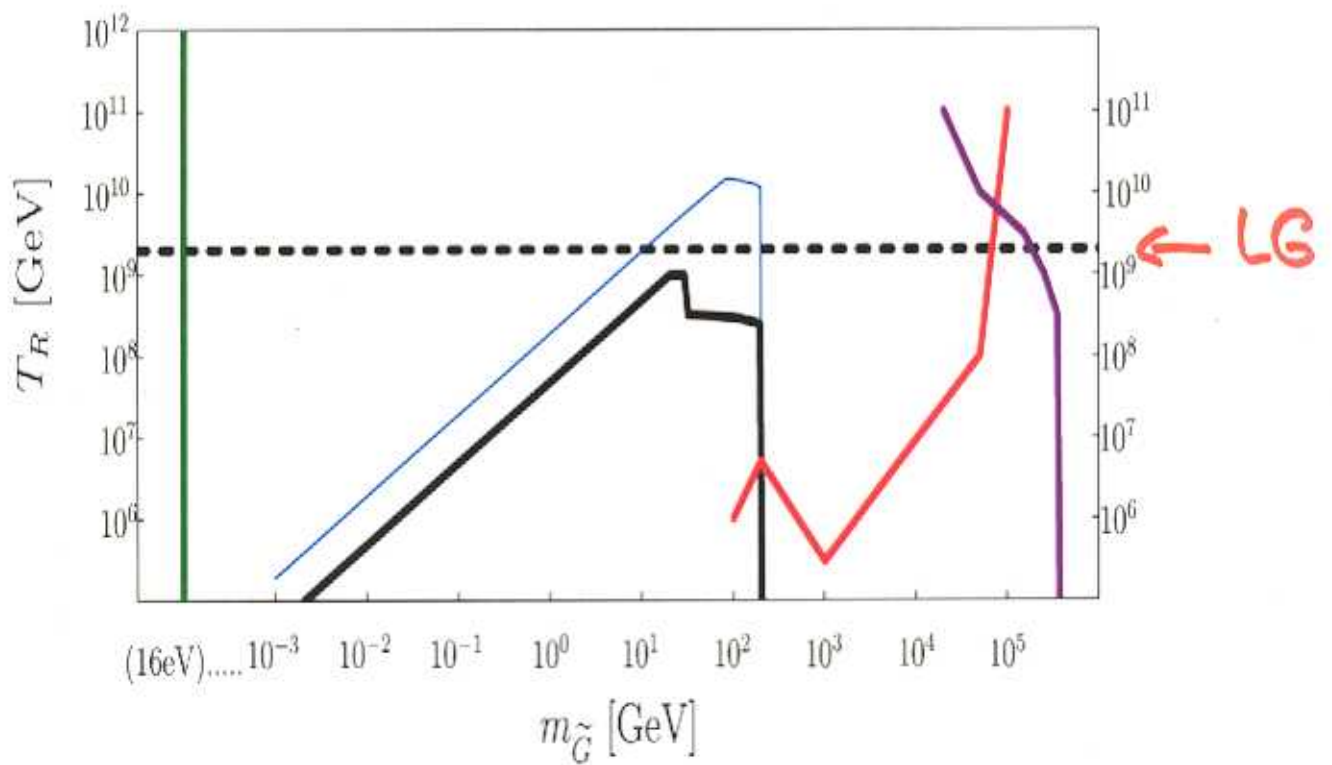
(i) BBN !

$$(ii) \Omega_{3/2}^{TH} h^2 + \Omega_{3/2}^{NT} h^2 \lesssim 0.1$$

Kawesaki, Kohri, Moroi hep-ph/0408426



Bound on reheating temperature



I

II

III

IV

V

—

$$m_{\tilde{g}} = 500 \text{ GeV}$$

—

$$m_{\tilde{g}} = 1 \text{ TeV}$$

\tilde{G} -summary : F

(i) \tilde{G} unstable

\tilde{U} : BBN , \tilde{D} : anomaly mediation,
 $\tilde{G} \rightarrow \pi_0 \gamma, \dots$ (Kitano, Ibe, Murayama '04)

typical \tilde{G} mass, $m_{3/2} \sim 1 \text{ TeV}$ worst!

only $m_{3/2} \sim 100 \text{ TeV}$ allowed!

(ii) \tilde{G} stable , $m_{3/2} \approx 100 \text{ GeV}$

$m_{\tilde{g}} = 1 \text{ TeV}$: only $m_{3/2} < 16 \text{ eV}$ (I)

allowed (LSS, Urd et al. '05)

$m_{\tilde{g}} = 500 \text{ GeV}$: $m_{3/2}$ around 50 GeV o.k. (iii)

(iii) the upper bound on T_R can be relaxed

- late entropy production

(Fujii, Yamagida ; WB, Hamaguchi, Ibe, Yan.)

- dynamical decrease of gauge couplings

for $T_R > T_* \approx (m_{3/2} M_p)^{1/2} < 10^{10} \text{ GeV}$

(WB, Hamaguchi, Ratz, '03)

$$\Omega_{3/2} h^2 \approx (0.05 - 0.2) \times \left(\frac{m_{\tilde{g}}}{1 \text{ TeV}} \right)^{3/2} \left(\frac{m_{3/2}}{1 \text{ GeV}} \right)^{1/4}$$

IV. Gravitino Dark Matter

Supersymmetry breaking on brane (\rightarrow gaugino mediation ...):

$$\mathcal{L}_D = \frac{1}{4g_D^2} \int d^2\theta W^\alpha W_\alpha + \text{h.c.} \\ + \delta^{(D-4)}(y-y_1) \frac{1}{4\Lambda} \int d^2\theta S W^\alpha W_\alpha + \text{h.c.}$$

$$\langle S \rangle = S_0 + \theta\theta F_S$$

$$\rightarrow m_{1/2} = \frac{g_{3/4}^2 F_S}{2\Lambda}, \quad m_{3/2} = \frac{1}{\sqrt{3}} \frac{F_S}{M_P} \quad (< m_{1/2}) \\ \text{(LSP)}$$

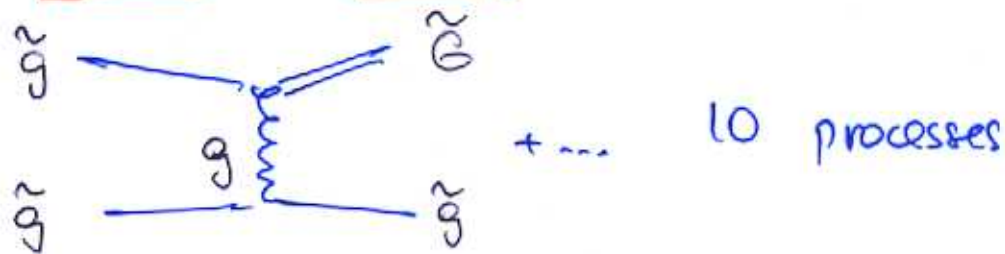
gauge coupling at finite temperature:

$$\frac{1}{g_4^2} + \frac{\Phi_T}{\Lambda} = g^2(\Phi_T), \quad \Phi = \text{Re} S$$

$$\text{for } T > T_* \sim \left(\frac{m_{3/2}^2 M_P}{2m_{\tilde{g}}^2} \right)^{1/2} :$$

$$g^2(T, \Phi_T) \approx \frac{g_4^2(T)}{1 + (T/T_*)^\alpha}, \quad \alpha = 1 \dots 4$$

gravitino production from thermal bath, after inflation, number density much below equilibrium density:



Boltzmann equations:

$$\frac{dn_{3/2}}{dt} + 3H n_{3/2} = C_{3/2}(T, \phi),$$

$$C_{3/2}(T, \phi) = -g^2(T, \phi) \frac{T^6}{M_p^2} \left(1 + \frac{\omega_{\tilde{g}}^2}{3\omega_{3/2}^2} \right)$$

$\text{spin} = \frac{3}{2}$
 $\text{spin} = \frac{1}{2}$

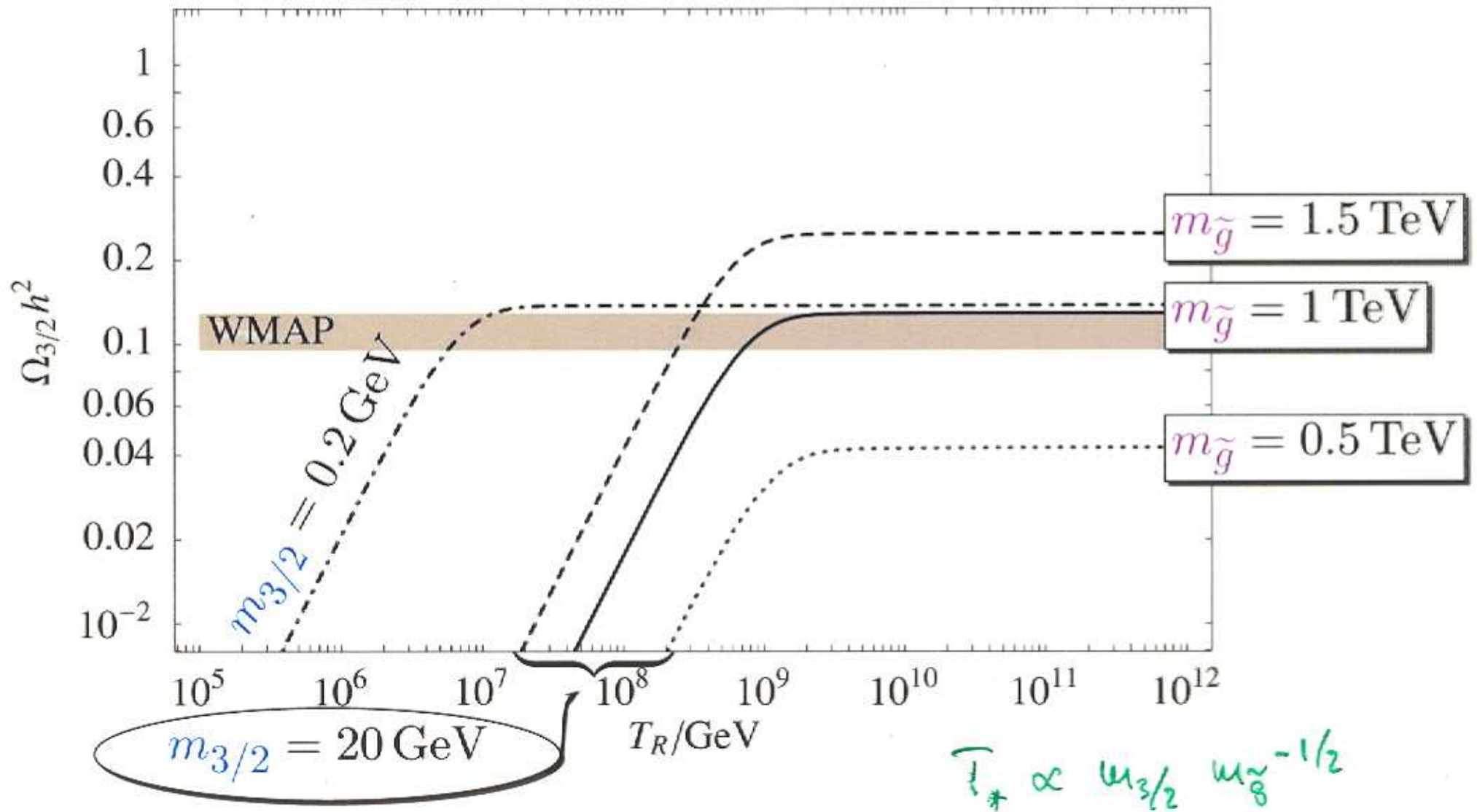
number density of gravitinos at present temperature T_0 for reheating T_R above T_* :

$$\frac{n_{3/2}}{s} \Big|_{T_0} = \frac{C_{3/2}(T_*, \phi)}{s(T_*) H(T_*)} \underbrace{P(d)}_{0.50 \dots 0.73, d=1 \dots 4}$$

temperatures above T_* don't contribute!

$$\frac{S_{3/2}}{s} \Big|_{T_0} = \dots \frac{T_*^2 \omega_{\tilde{g}}^2(T_*)}{M_p \omega_{3/2}} \propto \frac{\omega_{\tilde{g}}^{3/2}(T_*)}{M_p^{1/2}}$$

Contribution of gravitinos to Ωh^2



final result for $T_R > T_A$:

$$\Omega_{3/2} h^2 = 0.1 \times \left(\frac{m_{\tilde{g}}(1 \text{ TeV})}{1.0 \text{ TeV}} \right)^{3/2} \xi^{1/4} \hat{I}_a$$

$$\hat{I}_a = 0.5 \dots 2, \quad \xi = \mathcal{O}(1)$$

$$\text{WMAP: } \Omega_{\text{CDM}} h^2 = (\Omega_{\text{M}} - \Omega_{\text{b}}) h^2 = 0.113 \begin{matrix} +0.016 \\ -0.018 \end{matrix}$$

$$m_{\tilde{g}} = (0.5 \dots 2.0) \text{ TeV } \xi^{-1/6}$$

Ω_{CDM} depends only on $m_{\tilde{g}}$!

Simple picture of matter in the universe:

leptogenesis explains baryon density

$$\rightarrow T_R > \sim 10^{10} \text{ GeV}$$

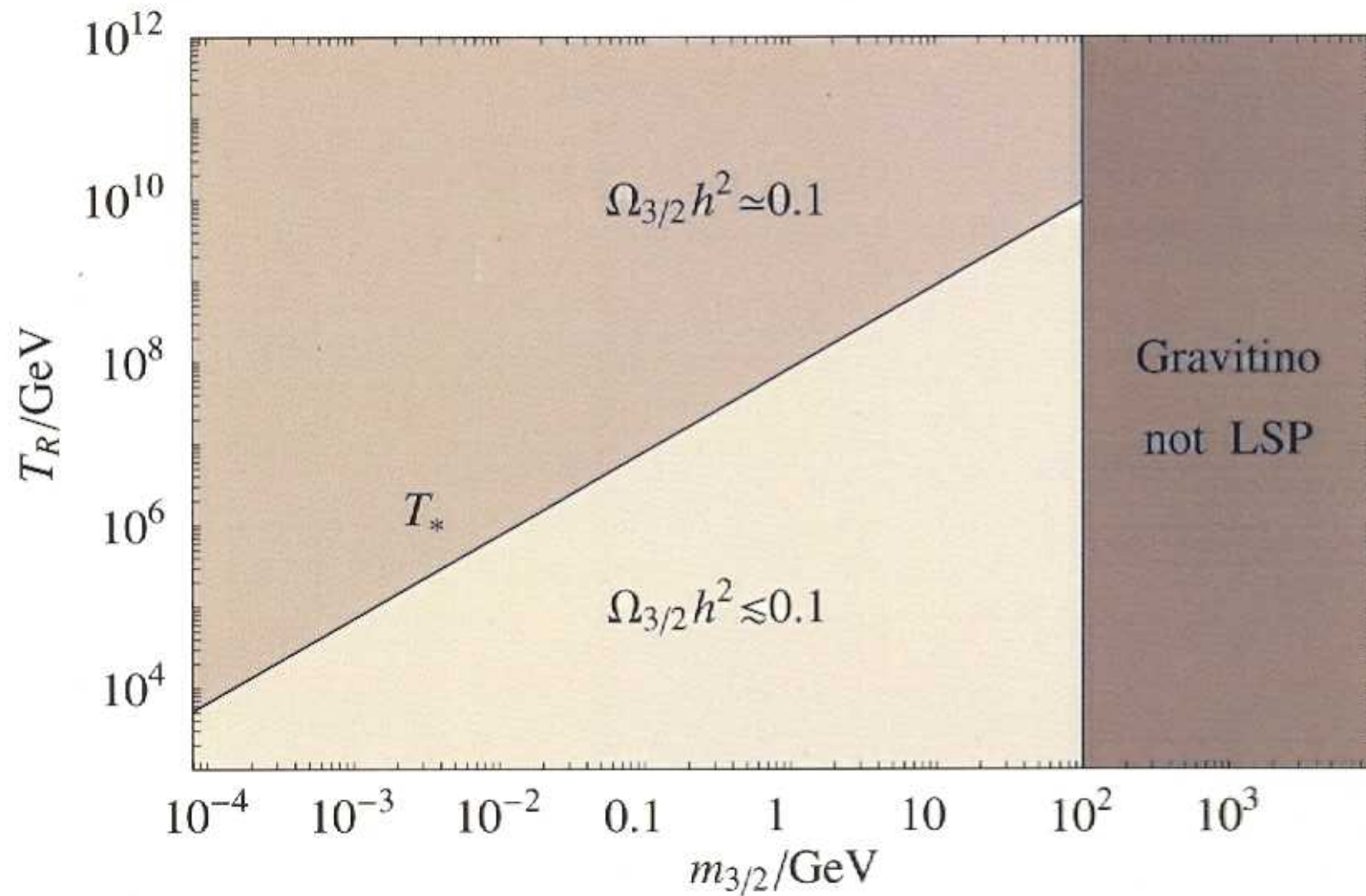
favours gravitino LSP, $m_{3/2} < 100 \text{ GeV}$

$$\rightarrow T_A^{\text{max}} < \sim 10^{10} \text{ GeV}$$

always reached; $\Omega_{3/2} h^2 \sim 0.1$ explained

for $m_{\tilde{g}} \sim 1 \text{ TeV} \rightarrow \text{LHC}$

Relic gravitino density



$$\xi_{3/2} = 1, \quad \omega_{\tilde{g}} = 1 \text{ TeV}$$