Perturbative QCD for the LHC

Nigel Glover

IPPP, Durham University

Highlights in Computational Quantum Field Theory

5th Vienna Central European Seminar

“Particle Physics and Quantum Field Theory”

Vienna, 28–30 November 2008

Present Status of QCD

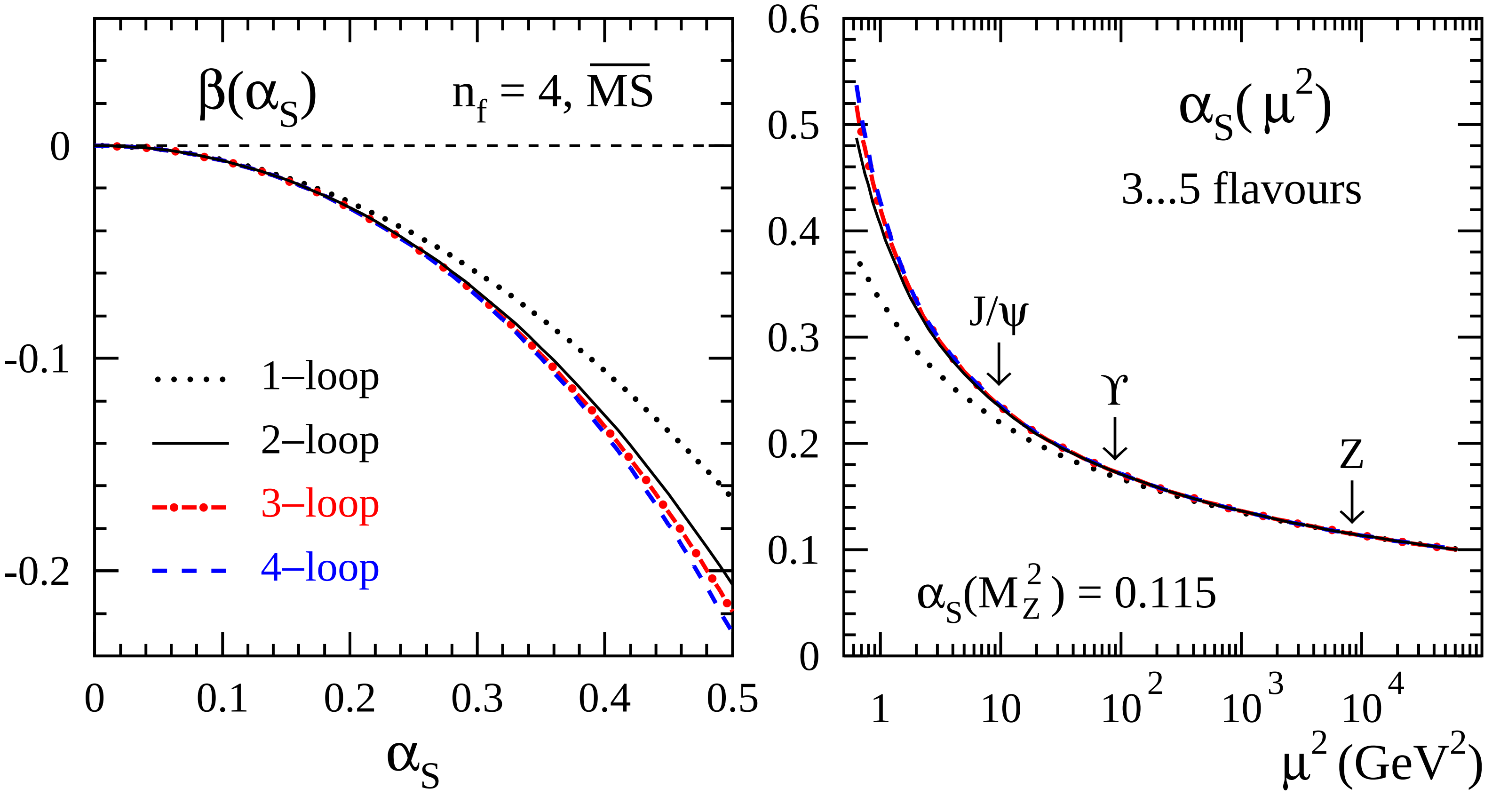
- ✓ Thanks to LEP, HERA and the TEVATRON
  QCD now firmly established theory of strong interactions
- ✓ We have gained a lot of confidence in comparing theoretical predictions with experimental data
- ✓ No major areas of discrepancies
- ✓ Now prepared to enter a new era of precision physics for QCD

The running coupling in perturbative QCD

$$d\alpha_s/d\ln\mu^2 = -\beta_0\,\alpha_s^2 - \beta_1\,\alpha_s^3 - \beta_2\,\alpha_s^4 - \beta_3\,\alpha_s^5 - \ldots$$

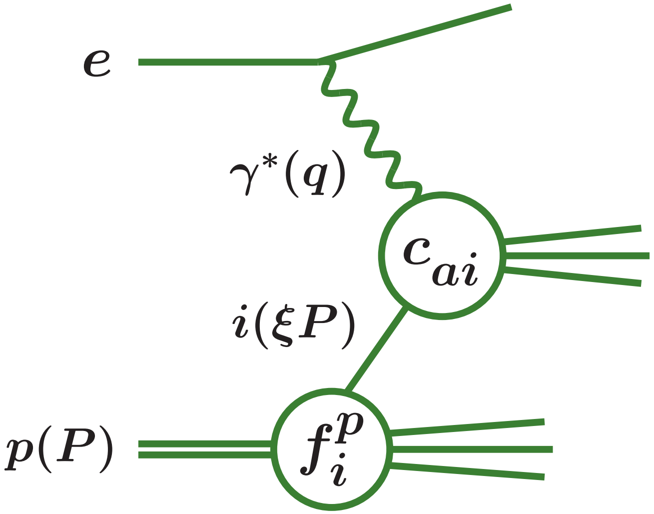
Four-loop coeff.:

van Ritbergen, Vermaseren, Larin; Czakon



Hard processes in perturbative QCD

Example: inclusive deep-inelastic scattering (DIS)



Kinematic variables

$$Q^2 = -q^2$$
$$x = Q^2/(2P\cdot q)$$

Lowest order : $x = \xi$

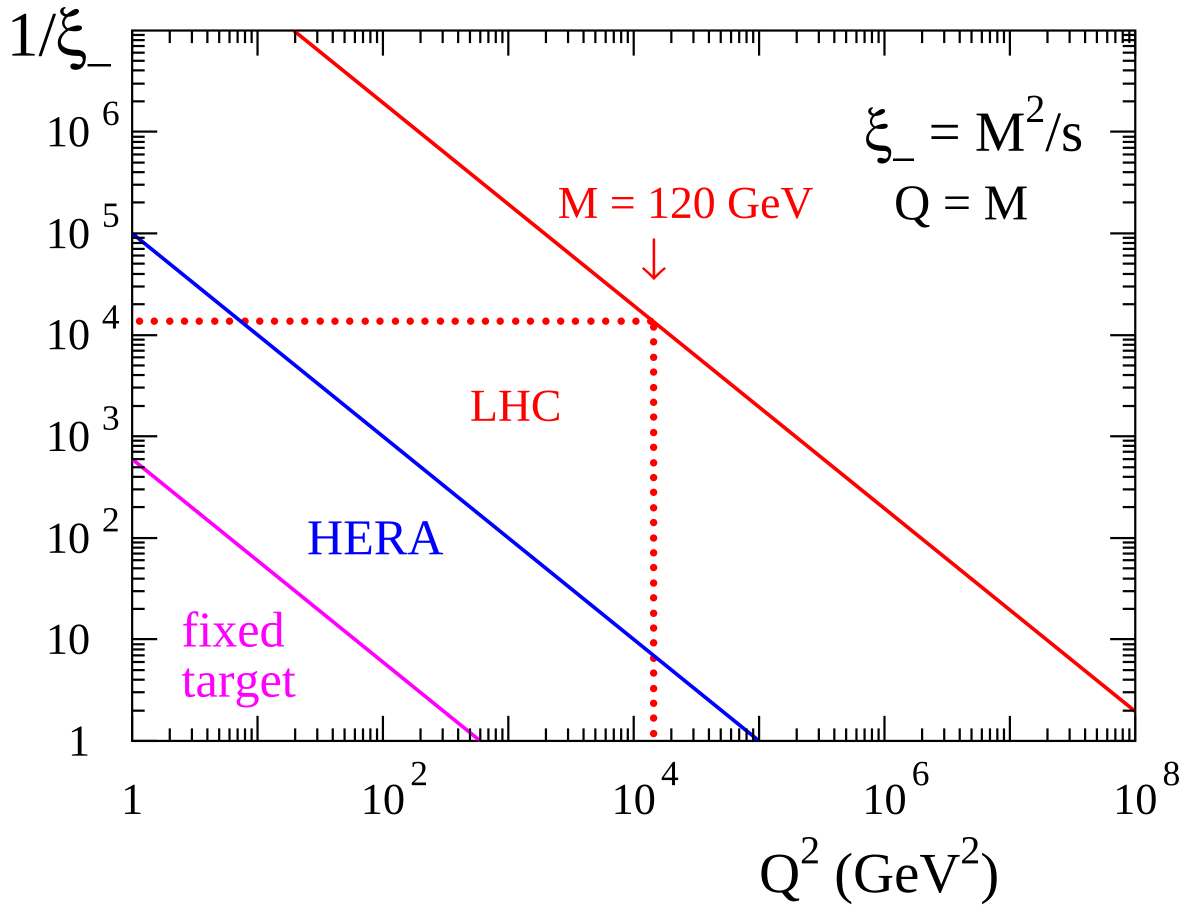
Structure functions F_a [up to $\mathcal{O}(1/Q^2)$]

$$F_a^p(x,Q^2) = \sum_i \left[c_{a,i}(\alpha_s(\mu^2),\mu^2/Q^2) \otimes f_i^p(\mu^2)\right](x)$$

Coefficient functions $c_{a,i}$, renormalization/factorization scale μ

Parton evolution from HERA to LHC

Kinematics: parton momenta $\xi_- < \xi < 1$ probed



HERA $\longrightarrow$ LHC:

Q^2 evolution across up to three orders of magnitude

Hard processes in perturbative QCD

Parton distributions f_i: evolution equations

$$\frac{d}{d\ln\mu^2} f_i(\xi,\mu^2) \;=\; \sum_k \left[P_{ik}(\alpha_s(\mu^2)) \otimes f_k(\mu^2)\right](\xi)$$

Initial conditions incalculable in pert. QCD.

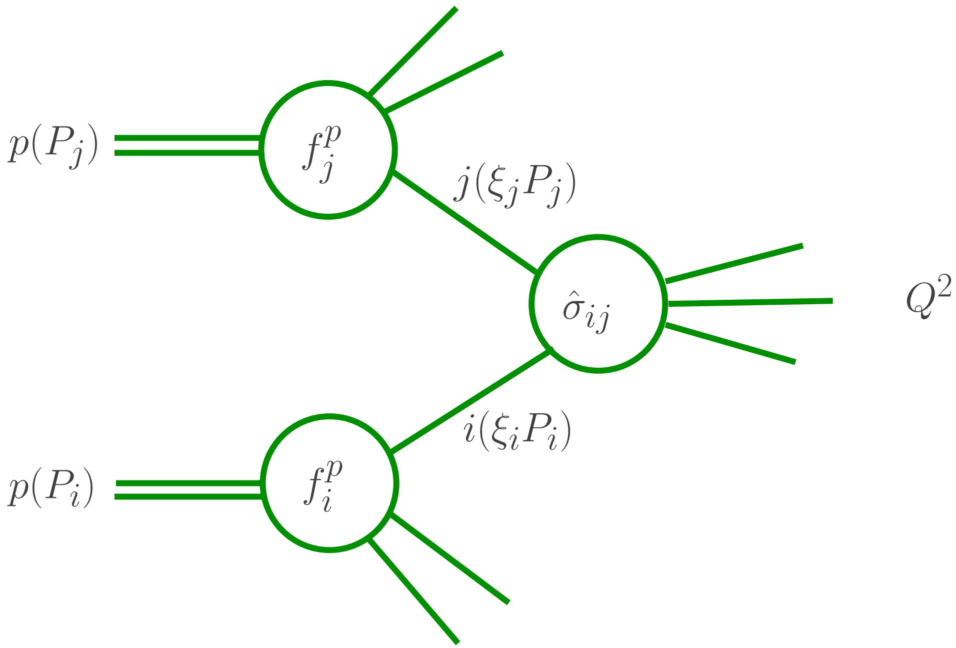
Splitting functions P, Coefficient functions c_a

$$\begin{aligned} P &= \alpha_s\, P^{(0)} + \alpha_s^2\, P^{(1)} + \alpha_s^3\, P^{(2)} + \ldots \\ c_a &= \alpha_s^{n_a}\left[c_a^{(0)} + \alpha_s\, c_a^{(1)} + \alpha_s^2\, c_a^{(2)} + \ldots\right] \end{aligned}$$

NLO: standard approximation
NNLO: new emerging standard

Moch, Vermaseren, Vogt

Hard processes in perturbative QCD



$$\sigma(Q^2) = \int \sum_{i,j} \left[d\hat{\sigma}_{ij}(\alpha_s(\mu_R), \mu_R^2/Q^2, \mu_F^2/Q^2) \otimes f_i^p(\mu_F) \otimes f_j^p(\mu_F) \right]$$

- ✓ partonic cross sections $d\hat{\sigma}_{ij}$
- ✓ running coupling $\alpha_s(\mu_R)$
- ✓ parton distributions $f_i(x, \mu_F)$
- ✓ renormalization/factorization scale μ_R, μ_F
- ✓ + parton shower + hadronisation model + underlying event + ...

The challenge

- ✓ Everything at the LHC (signals, backgrounds, luminosity measurement) involves QCD
- ✓ Strong coupling is not small: $\alpha_s(M_Z) \sim 0.12$ and running is important
  - ⇒ events have high multiplicity of hard partons
  - ⇒ each hard parton fragments into a cluster of collimated particles jet
  - ⇒ higher order perturbative corrections can be large
  - ⇒ theoretical uncertainties can be large
- ✓ Processes can involve multiple energy scales: e.g. p_T^W and M_W
  - ⇒ may need resummation of large logarithms
- ✓ Parton/hadron transition introduces further issues, but for suitable (infrared safe) observables these effects can be minimised
  - ⇒ importance of infrared safe jet definition
  - ⇒ accurate modelling of underlying event, hadronisation, ...

What is covered in this talk

Will focus on status of fixed order parton-level predictions

- ✓ Systematic to higher order/higher multiplicity in perturbation theory
- ✓ Appropriate for hard well separated final states
- ✓ Lead to a systematic reduction in renormalisation/factorisation scale uncertainties
- ✓ Many recent theoretical developments and new calculations/numerical programmes available

caveat Parton-level, relies on matching to experimental observables e.g. merging with parton showers and event generators, etc

CKKW, MLM, MCNLO, POWHEG

see talk by Seymour

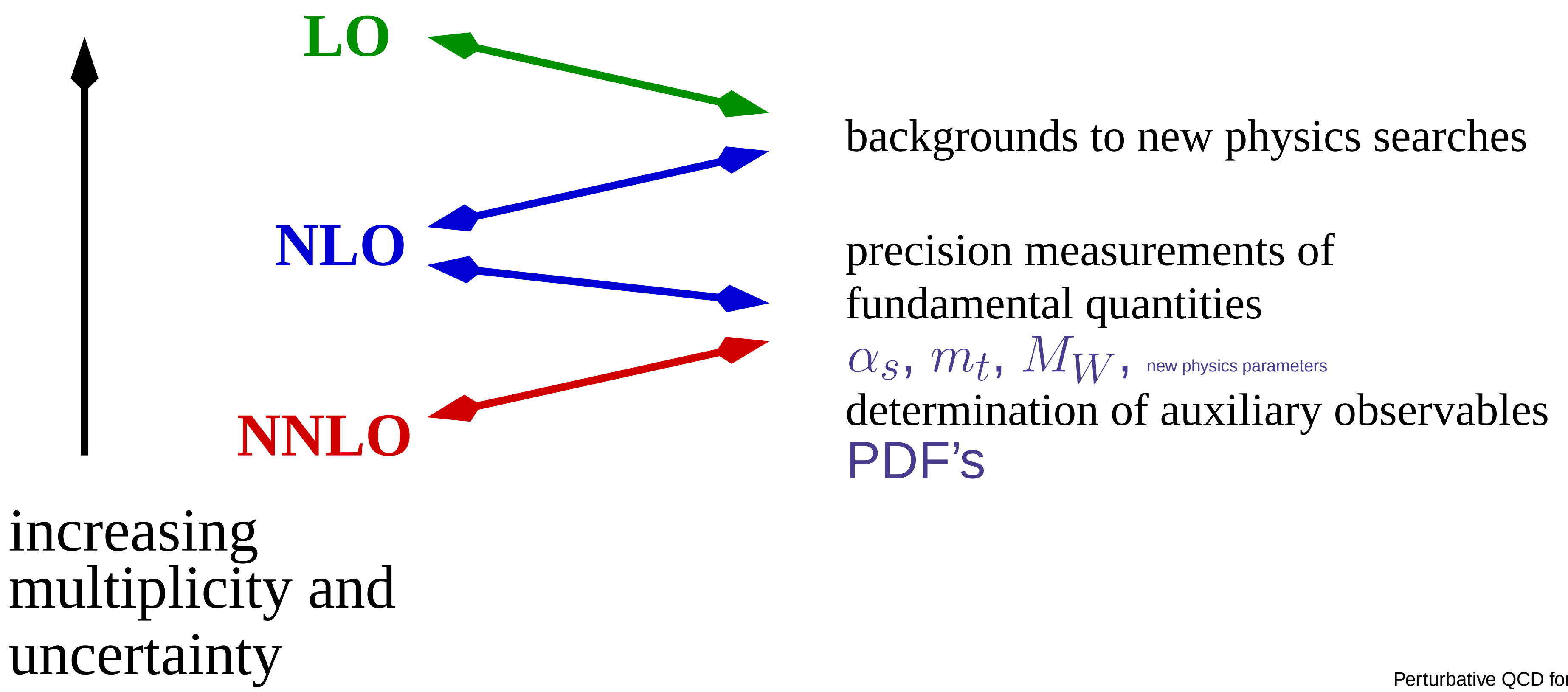
- ✗ No time for many important topics;
  - ✗ parton distributions
  - ✗ soft gluon resummation
  - ✗ studies of jet definitions; fast k_T algorithm, infrared safe cone algorithms,...

Salam et al

Matching onto Physics Goals

Twin Goals:

1. Identification and study of New Physics
2. Precision measurements (e.g. α_s, PDF's) leading to improved theoretical predictions



State of the Art - at a glance

Relative Order	$2 \to 1$	$2 \to 2$	$2 \to 3$	$2 \to 4$	$2 \to 5$	$2 \to 6$
1	LO					
α_s	NLO	LO				
α_s^2	NNLO	NLO	LO			
α_s^3		NNLO	NLO	LO		
α_s^4				NLO	LO	
α_s^5					NLO	LO

LO Well under control, even for multiparticle final states

NLO Well understood for $2 \to 1$ and $2 \to 2$

NLO Many new $2 \to 3$ calculations

NLO Very first $2 \to 4$ LHC cross section this year

NNLO Recent breakthroughs for inclusive and exclusive $2 \to 1$

NNLO Recent landmark calculation of NNLO splitting functions

Moch, Vermaseren, Vogt

NNLO Still waiting for $2 \to 2$

Leading order

Many available programs for automatic evaluation of LO cross sections

- ✓ Feynman diagrams: matrix elements automatically generated up to $2 \to 6$
  MADGRAPH, COMPHEP, GRACE, ...
- ✓ Off-shell recursion relations:

  Berends, Giele; Caravaglios, Moretti

  matrix elements automatically generated up to $2 \to 8$ or more
  HELAC, AMEGIC++/COMIX, ALPHA, ...
- ✓ (Twistor inspired) On-shell recursion relations:

  Cachazo, Svrcek, Witten; Britto, Cachazo, Feng, Witten, AMEGIC++/COMIX; Dinsdale, Ternick, Weinzierl

- ✓ plus automatic integration over phase space
  HELAC/PHEGAS, MADGRAPH/MADEVENT, SHERPA/AMEGIC++, ALPHA/ALPGEN, ...
- ✓ very good for estimating importance of various processes in different models - properly populate phase space with multiple hard objects
- ✓ able to interface with parton showers CKKW in SHERPA, MLM in ALPGEN, ...

Comparison of algorithms

- ✓ On-shell recursion relations (CSW, BCF) yield compact analytic results
- ✓ Numerical implementations show that off-shell recursion (BG) is faster

Final state	BG		BCF		CSW	
	CO	CD	CO	CD	CO	CD
2g	0.24	0.28	0.28	0.33	0.31	0.26
3g	0.45	0.48	0.42	0.51	0.57	0.55
4g	1.20	1.04	0.84	1.32	1.63	1.75
5g	3.78	2.69	2.59	7.26	5.95	5.96
6g	14.20	7.19	11.90	59.10	27.80	30.60
7g	58.50	23.70	73.60	646	146	195
8g	276	82.10	597	8690	919	1890
9g	1450	270	5900	127000	6310	29700
10g	7960	864	64000		48900	

Duhr, Hoche, Maltoni

Example at LO

Multi-jet production at the LHC using HELAC/PHEGAS

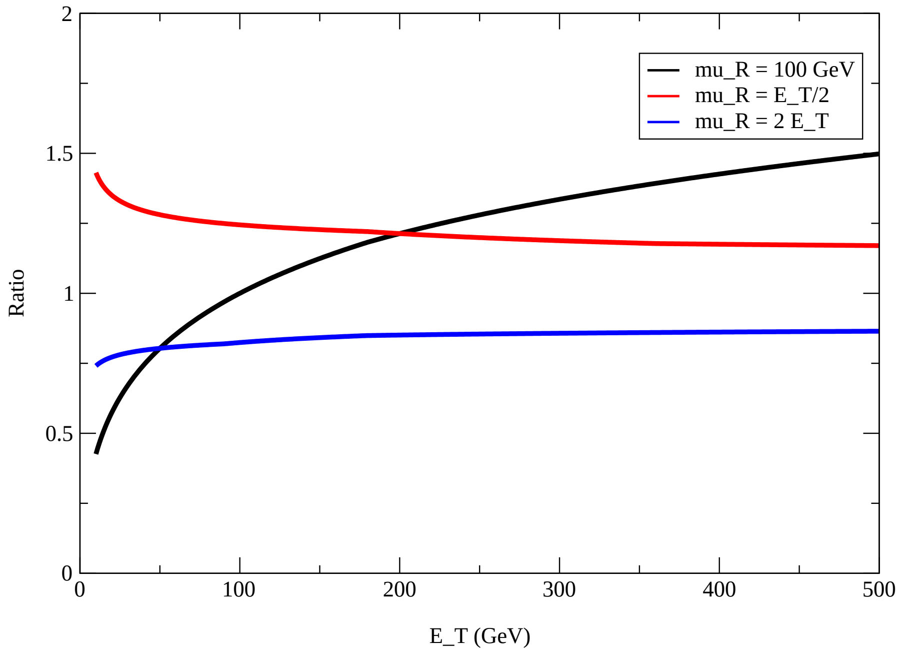
Draggiotis, Kleiss, Papadopoloulos

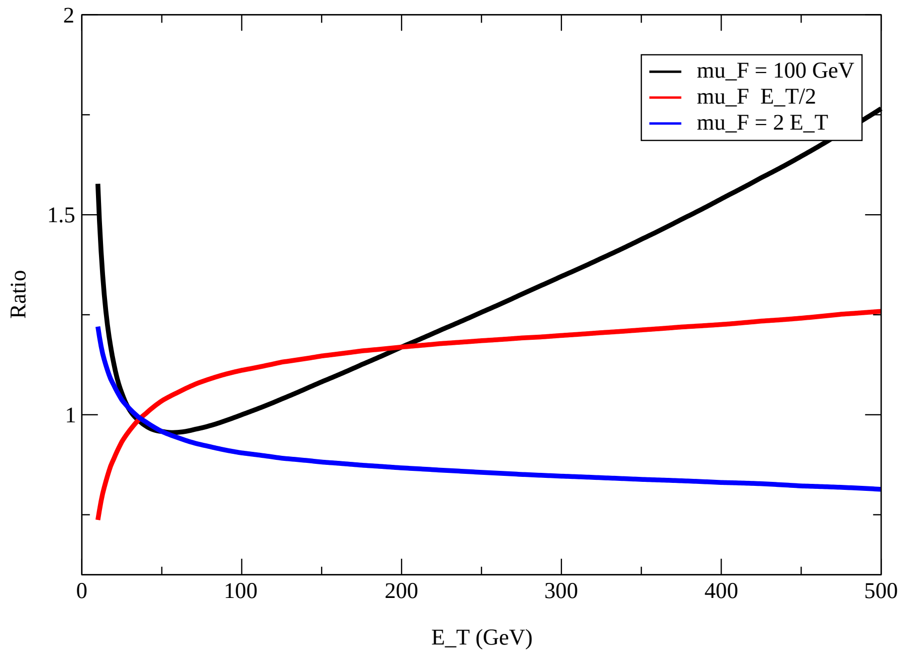
# of jets	2	3	4	5	6	7	8
# of dist.processes	10	14	28	36	64	78	130
total # of processes	126	206	621	861	1862	2326	4342
$\sigma(nb)$	-	91.41	6.54	0.458	0.030	0.0022	0.00021
$\%$ Gluonic	-	45.7	39.2	35.7	35.1	33.8	26.6

- ✓ For each final state, there are many distinct contributing processes e.g. $gg \to gg$, $gg \to q\bar{q}$, $q\bar{q} \to gg$, $qg \to qg$, $q\bar{q} \to Q\bar{Q}$, $qQ \to qQ$ etc
- ✓ Assigning different quark flavours gives even more
- ✓ Bookkeeping, phase space generation and evaluation done automatically
- ✓ ALPGEN and SHERPA are also very fast for multiparticle SM processes
- ✓ MADGRAPH slower, but adapted for other models, effective H, MSSM, 2HDM, ...

Limitations of LO

Very large uncertainty for multiparticle final states



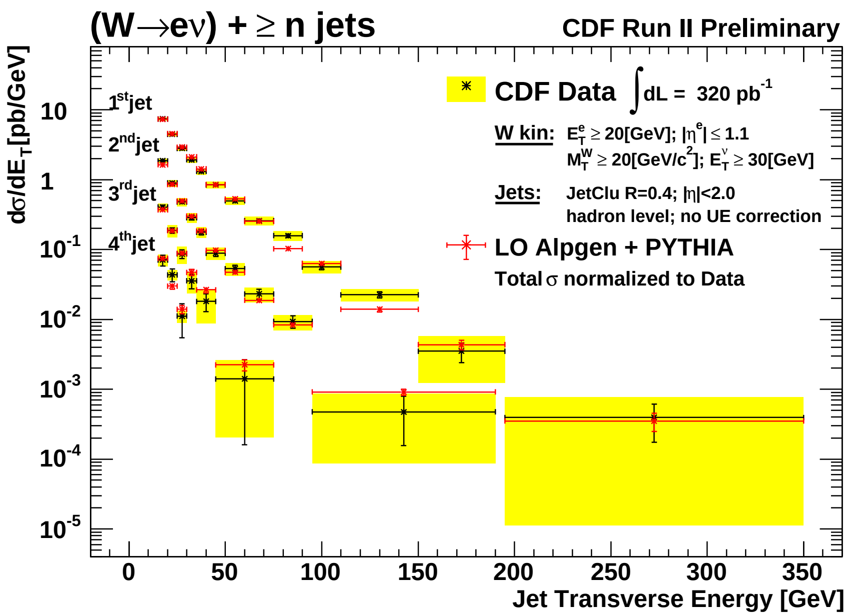


scale uncertainty on α_s^2

parton luminosity uncertainty

- ✓ New channels open up at higher orders qg + large gluon PDF
- ✓ Increased phase space
- ✓ Large π^2 coefficients in s-channel $\Rightarrow$ large NLO corrections 30% - 100%

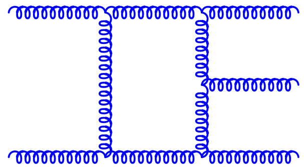
W + Jets at CDF Run II with 320 pb^{-1}



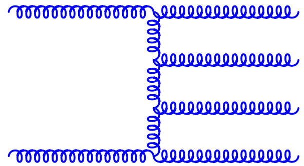
cross sections for the leading jet in $W+ \geq 1$ jet events, second jet in $W+ \geq 2$ jets events, etc

- ✓ ALPGEN+PYTHIA merged LO+PS prediction normalised to the inclusive cross section for each jet multiplicity
- ✓ Excellent qualitative agreement

Anatomy of a NLO calculation



✓ one-loop $2 \rightarrow 3$ process
looks like 3 jets in final state



✓ tree-level $2 \rightarrow 4$ process
looks like 3 or 4 jets in final state

✓ plus method for combining the infrared divergent parts - dipole subtraction

Catani, Seymour; Dittmaier, Trocsanyi, Weinzierl, Phaf

✓ automated dipole subtraction

Gleisberg, Krauss (SHERPA); Hasegawa, Moch, Uwer; Frederix, Gehrmann, Greiner (MadDipole); Seymour, Tevlin – see talk by Hasegawa

Bottleneck: one-loop matrix elements

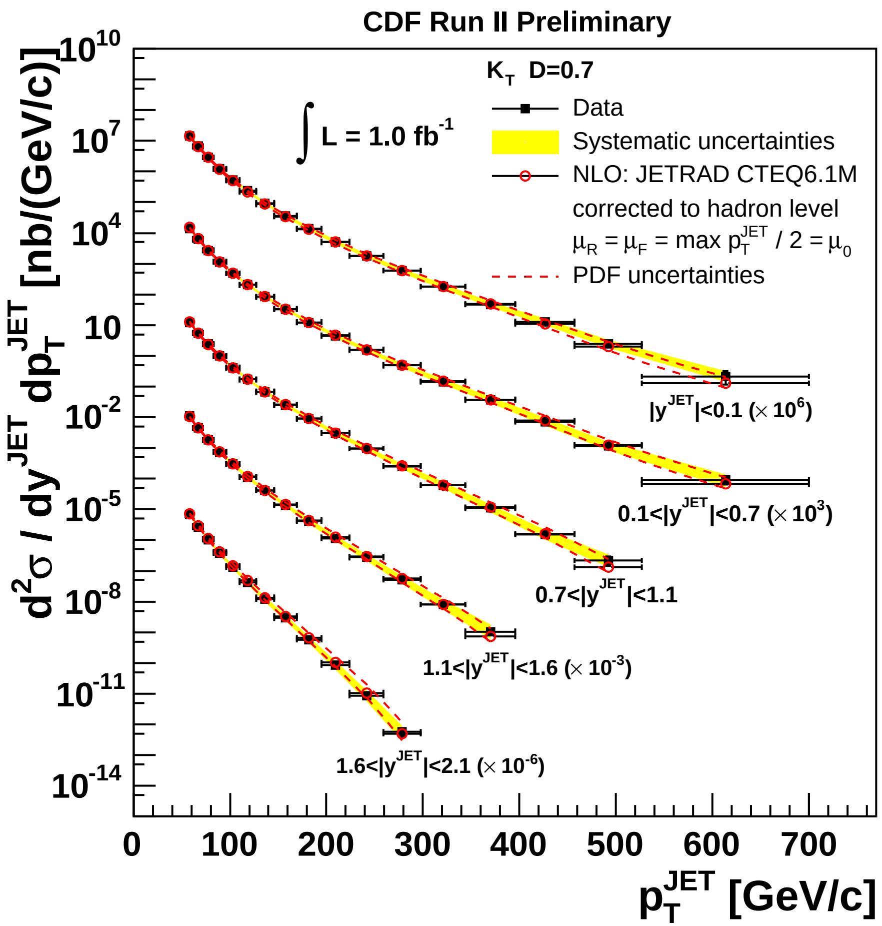
Availability of NLO calculations

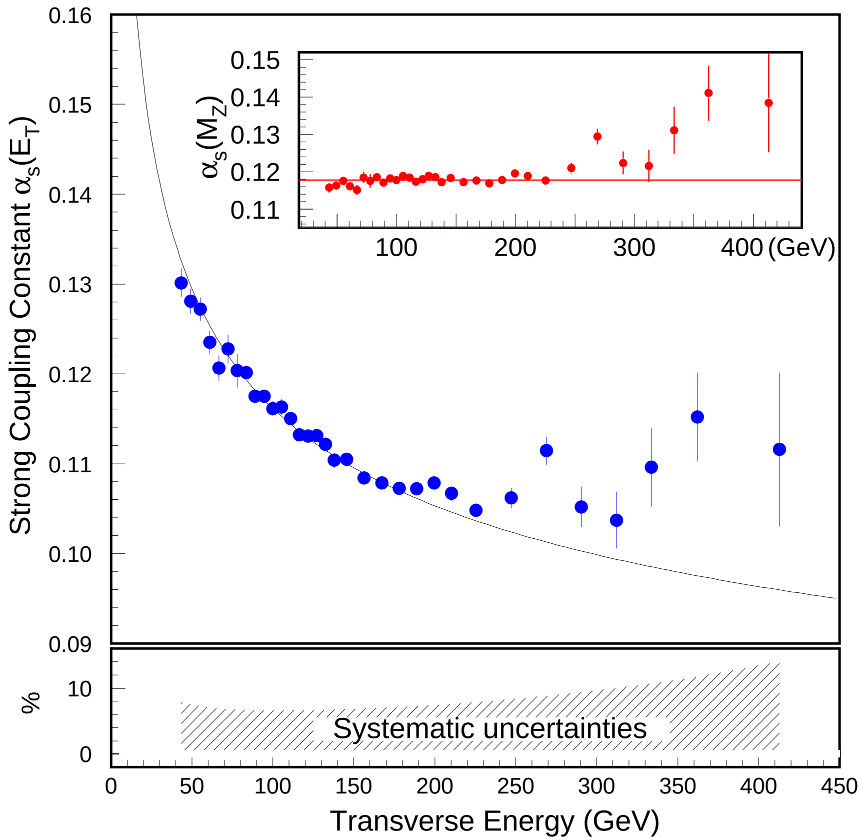
- ✓ $2 \to 2$ processes
  - ✓ parton level integrators available for all $2 \to 2$ Standard Model and MSSM processes for some time
  - ✓ extensively used at LEP, TEVATRON and HERA
    EVENT, JETRAD, MCFM, DISENT, DIPHOX, HQQB, NLOJET++, VBFNLO etc
  - ✓ can be matched with parton shower MC@NLO, POWHEG – Frixione, Webber; Nason, Oleari, Ridolfi; Krämer, Soper see talk by Seymour
- ✓ $2 \to 3$ processes
  - ✓ many $2 \to 3$ processes now available at NLO
    e.g. backgrounds $pp \to 3$ jets, $V + 2$ jets, $\gamma\gamma$ + jet, $V + b\bar{b}$, VV + jet, $t\bar{t}+$ jet
    as well as signals $pp \to t\bar{t}H$, $b\bar{b}H$, $H + 2$ jets, HHH, $t\bar{t}$+jet

  http://www.cedar.ac.uk/hepcode

⇒ First $2 \to 4$ LHC cross section recently computed

Inclusive Jet Production using the Kt Algorithm





Single jet inclusive differential cross section in different rapidity slices

✓ Described by NLO QCD

✓ Excellent quantitative agreement $\Longrightarrow$ CDF Run I α_s measurement

LHC priority NLO wish list, Les Houches 2005/7*

process	background	status
$pp \to VV + 1$ jet	WBF $H \to VV$	WWj (07)
$pp \to t\bar{t} + b\bar{b}$	$t\bar{t}H$	$q\bar{q} \to t\bar{t}b\bar{b}$ (08)
$pp \to t\bar{t} + 2$ jets	$t\bar{t}H$	$t\bar{t}j$ (07)
$pp \to VV + b\bar{b}$	WBF $H \to VV$, $t\bar{t}H$, NP	
$pp \to VV + 2$ jets	WBF $H \to VV$	WBF $pp \to VVjj$ (07)
$pp \to V + 3$ jets	NP	
$pp \to VVV$	SUSY trilepton	ZZZ (07), WWZ (07), WWW (08), ZZW (08)
$pp \to b\bar{b}b\bar{b}$*	Higgs and NP	

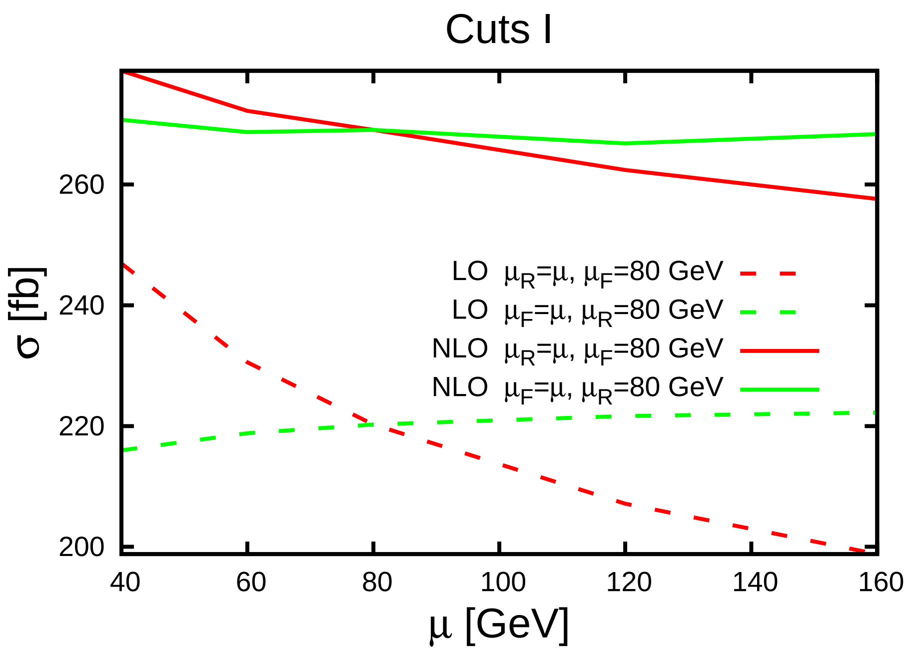
- ✓ $pp \to H + 2$ jets via gluon fusion — Campbell, Ellis, Zanderighi, hep-ph/0608194
- ✓ $pp \to H + 2$ jets via WBF, electroweak and QCD corrections — Ciccolini, Denner, Dittmaier, arXiv/0710.4749
- ✓ $pp \to H + 3$ jets via WBF, — Figy, Hankele, Zeppenfeld, arXiv/0710.5621
- ✓ . . . .

Vector boson pair plus jet

QCD corrections to $pp \to W^+W^-j + X$

Dittmaier, Kallweit, Uwer, arXiv/0710.1577; Campbell, Ellis, Zanderighi, arXiv/0710.1832

✓ Background to Higgs in both WBF, GF channels - $H \to W^+W^-$ with one jet missed, or Higgs recoiling against jet



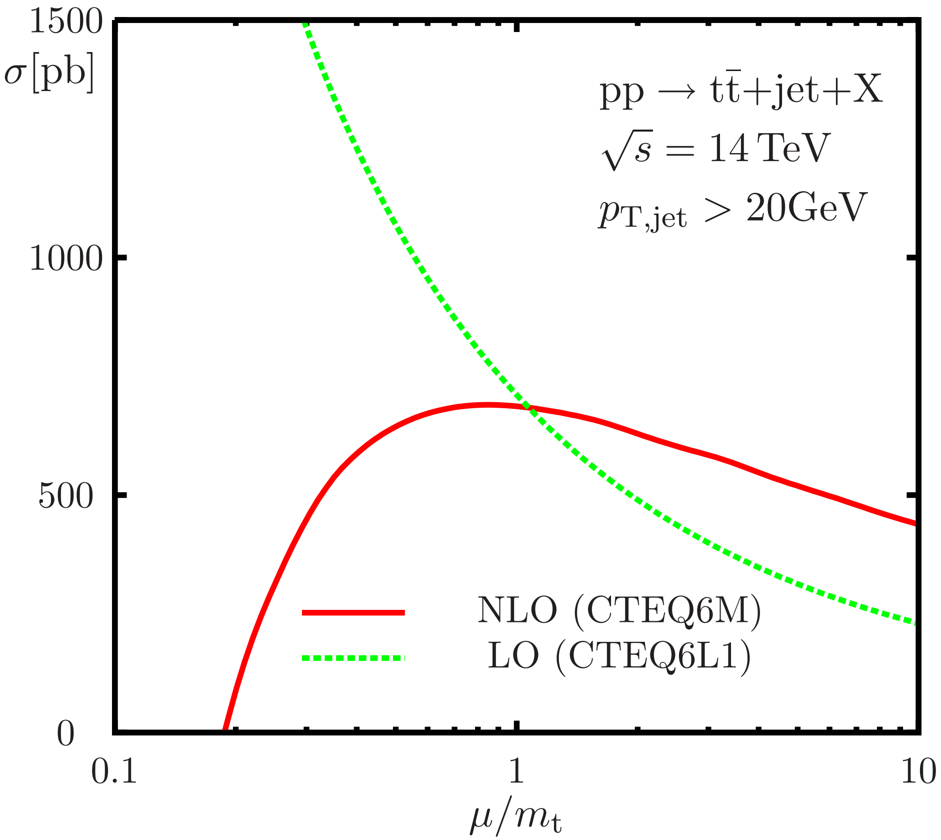
- ✓ For inclusive cuts, NLO increases cross section by about 25%
- ✓ Factorisation scale uncertainty small, renormalisation scale uncertainty reduced by $\sim$ 50%
- ✓ Shapes of NLO inclusive distributions very similar to LO
- ✓ For WBF cuts, with one or both jets forward, WWj is one of dominant backgrounds
  NLO increased by $\sim$ 70% cf LO

Top pair plus jet

QCD corrections to $pp \to t\bar{t}j + X$

Dittmaier, Uwer, Weinzierl hep-ph/0703120

- ✓ Background to Higgs in WBF, $t\bar{t}H$ channels
- ✓ measurement of t properties



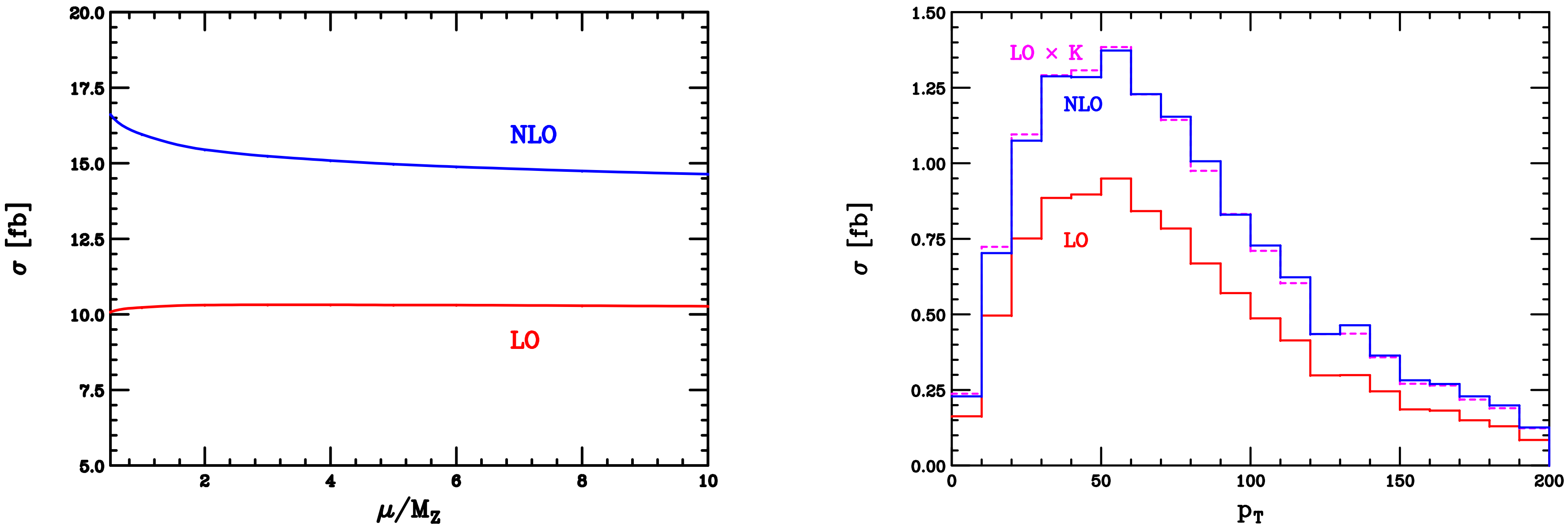
- ✓ Residual scale dependence reduced
- ✓ NLO corrections essentially eliminate forward-backward charge asymmetry at Tevatron

Triple Vector Boson Production

QCD corrections to $pp \to ZZZ + X$

Lazopoulos, Melnikov, Petriello, hep-ph/0703273

✓ Background to various SUSY tri-lepton signatures, gauge boson coupling measurments,



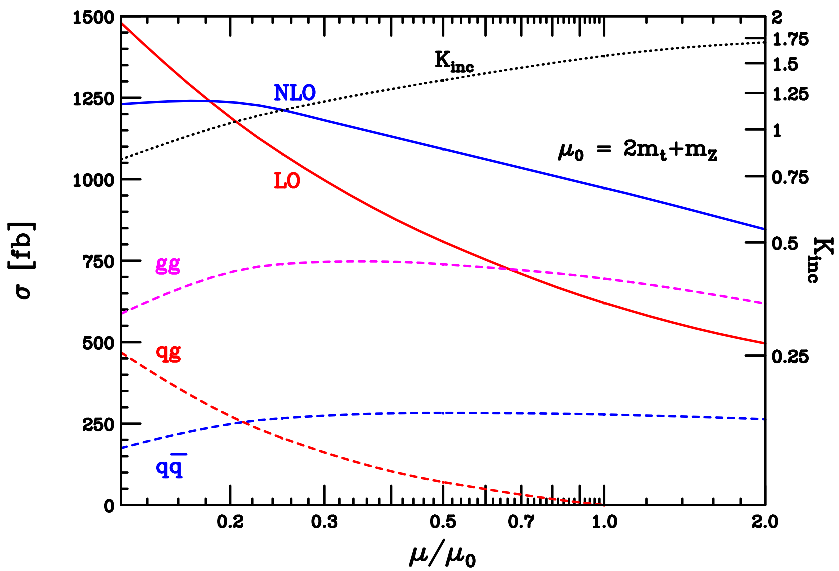
✓ Large, 50% corrections not seen by LO scale variation! 15% shift from pdfs, 35% shift from π^2 terms

Top pair plus Z Production

QCD corrections to $pp \to t\bar{t}Z + X$

Lazopoulos, Melnikov, Petriello, arXiv/0709.4044,
Lazopoulos, McElmurry, Melnikov, Petriello, arXiv/0804.2220

- ✓ Background to various SUSY tri-lepton signatures, gauge boson coupling measurments,
- ✓ Fully numerical calculation - using sector decomposition and contour deformation
- ✓ For reasonable choices of μ, corrections as large as 75%

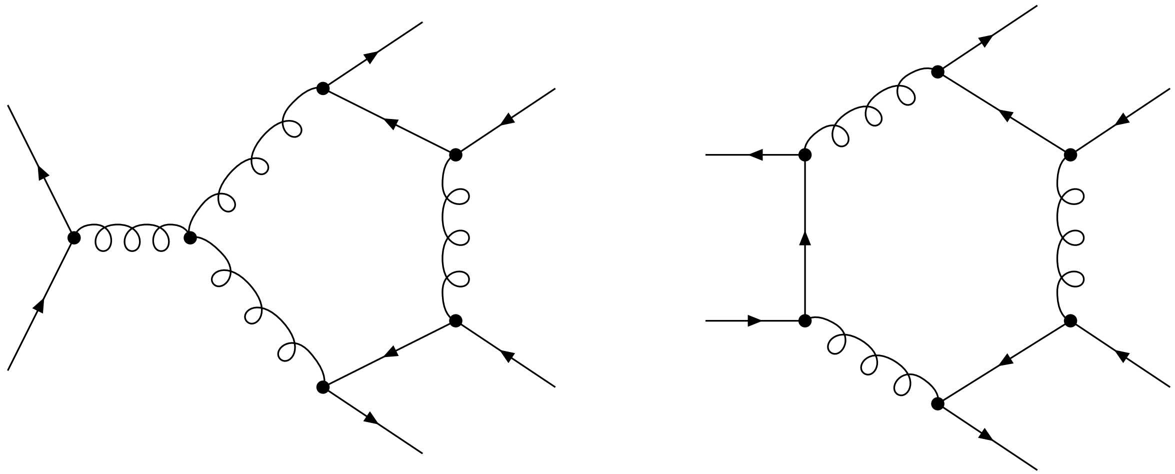


Top pair plus bottom pair Production

QCD corrections to $q\bar{q} \to t\bar{t}b\bar{b} + X$

Bredenstein, Denner, Dittmaier, Pozzorini, arXiv/0807.1248

✓ Background to the Higgs signal in $t\bar{t}H$ production where the Higgs decays into a bottom pair



✓ First successful demonstration of Feynman diagrammatic evaluation of $2 \to 4$ process at LHC

✓ Dominant $gg \to t\bar{t}b\bar{b} + X$ process underway

The one-loop problem

Any one-loop integral can be written as

$$\begin{aligned}\mathcal{M} &= \sum d(D)\mathrm{boxes(D)} \\ &+ \sum c(D)\mathrm{triangles(D)} \\ &+ \sum b(D)\mathrm{bubbles(D)} \\ &+ \sum a(D)\mathrm{tadpoles(D)}\end{aligned}$$

- ✓ most of the scalar loop integrals boxes etc are known analytically around $D = 4$
- ✓ only problem is to compute the D-dimensional coefficients $a(D)$ etc.

The only problem is complexity - the number of terms generated is too large to deal with, even with computer algebra systems, and there can be very large cancellations.

The one-loop problem - continued

Lots of ideas and strategies

- ✓ Improved tensor reduction: Denner, Dittmaier; Binoth, Guillet, Heinrich, Pilon, Reiter, Schubert, . . . – see talk by Reiter
- ✓ Numerical evaluation of integral recursion relations Giele, Ellis, Zanderighi
- ✓ 4-d Unitarity and cut constructibility Bern, Dixon, Dunbar, Kosower; Britto, Cachazo, Feng; . . .
- ✓ D-dimensional unitarity Anastasiou, Britto, Cachazo, Feng, Kunszt, Mastrolia
- ✓ Numerical loop integration: accuracy only has to match real emission contribution Nagy, Soper, hep-ph/0610028
  Sector decomposition plus contour deformation automated by Anastasiou, Beerli, Daleo, hep-ph/0703282
- ✓ Algebraic reduction of the integrand
  Ossola, Papadopoulos, Pittau, hep-ph/0609007; Ellis, Giele, Kunszt, arXiv/0708.2398
- ✓ . . . . . .

Testing ground: Six-photon amplitude hep-ph/0610028, hep-ph/0703311 hep-ph/0704.1271

One-loop six gluon amplitude

✓ Analytic computation
Bedford, Berger, Bern, Bidder, Bjerrum-Bohr, Brandhuber, Britto, Buchbinder, Cachazo, Dixon, Dunbar, Feng, Forde, Kosower, Mastrolia, Perkins, Spence, Travaglini, Xiao, Yang, Zhu

Amplitude	$\mathcal{N} = 4$	$\mathcal{N} = 1$	$\mathcal{N} = 0$ (cut)	$\mathcal{N} = 0$ (rat)
$--++++$	BDDK (94)	BDDK (94)	BDDK (94)	BDK (94)
$-+-+++$	BDDK (94)	BDDK (94)	BBST (04)	BBDFK (06), XYZ (06)
$-++-++$	BDDK (94)	BDDK (94)	BBST (04)	BBDFK (06), XYZ (06)
$---+++$	BDDK (94)	BDDK (94)	BBDI (05), BFM (06)	BBDFK (06), XYZ (06)
$--+-++$	BDDK (94)	BBDP (05), BBCF (05)	BFM (06)	XYZ (06)
$-+-+-+$	BDDK (94)	BBDP (05), BBCF (05)	BFM (06)	XYZ (06)

✓ Numerical evaluation via recursion Ellis, Giele, Zanderighi (06)

✓ Numerical evaluation based on unitarity Ellis, Giele, Kunszt (07)

The multi-leg one-loop problem - 1

For theories with massless internal particles, it might be easier to compute

$$\begin{aligned}\mathcal{M} &= \sum d(4)\mathrm{boxes}(\mathrm{D}) \\ &+ \sum c(4)\mathrm{triangles}(\mathrm{D}) \\ &+ \sum b(4)\mathrm{bubbles}(\mathrm{D}) \\ &+ \mathrm{R}\end{aligned}$$

where the coefficients are now 4-dimensional and R is a rational (non-logarithmic) term

- ✓ Coefficients of loop integrals obtained with generalised unitarity
- ✓ Rational parts obtained using on-shell recursion relation.

Berger, Bern, Dixon, Forde, Kosower

- ✓ Numerical implementation in `BlackHat`

Berger, Bern, Dixon, Febres Cordero, Forde, Ita, Kosower, Maitre

The multi-leg one-loop problem - 2

Alternatively

$$
\begin{aligned}
\mathcal{M} \quad = & \qquad\qquad\qquad\qquad\qquad\quad + \sum e'()\mathrm{pentagons}(\mathrm{D}+2) \\
& + \sum d()\mathrm{boxes}(\mathrm{D}) \quad +d'()(D-4)\mathrm{boxes}(\mathrm{D}+2) \\
& \qquad\qquad\qquad\qquad\quad +d''()(D-4)(D-2)\mathrm{boxes}(\mathrm{D}+4) \\
& + \sum c()\mathrm{triangles}(\mathrm{D}) \quad +c'()(D-4)\mathrm{triangles}(\mathrm{D}+2) \\
& + \sum b()\mathrm{bubbles}(\mathrm{D}) \quad +b'()(D-4)\mathrm{bubbles}(\mathrm{D}+2)
\end{aligned}
$$

where the coefficients don't depend on the dimension.

Giele, Kunszt, Melnikov

✓ To probe the higher dimension terms, one can employ unitarity in higher (integer) dimension

Ellis, Giele, Kunszt, Melnikov

✓ Numerical implementation in `ROCKET`

Giele, Zanderighi

Numerical one-loop evaluation: BlackHat v Rocket

BlackHat

- ✓ based on $D = 4$ unitarity and on-shell recursion
- ✓ up to 8 gluon amplitudes numerically

Berger et al, arXiv:0803.4180

- ✓ leading colour $Vq\bar{q}ggg$ numerically

Berger et al, arXiv:0808.0941

Rocket

- ✓ based on OPP and D_s dimension unitarity
- ✓ off-shell recursion for tree-input
- ✓ up to 20 gluon amplitudes numerically

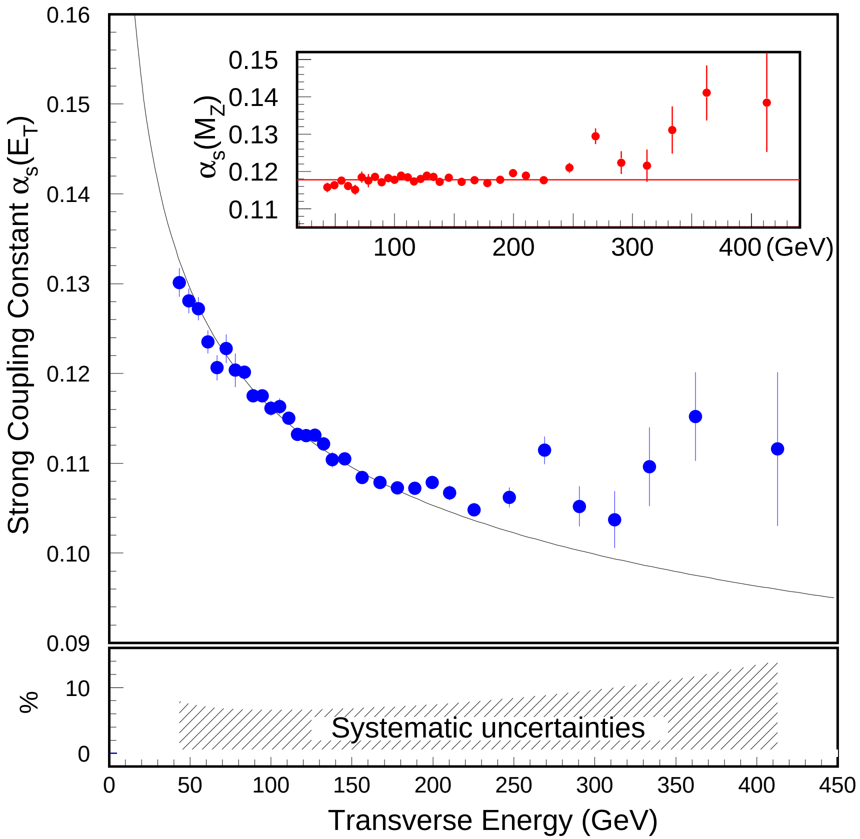
Giele, Zanderighi, arXiv:0805.2152

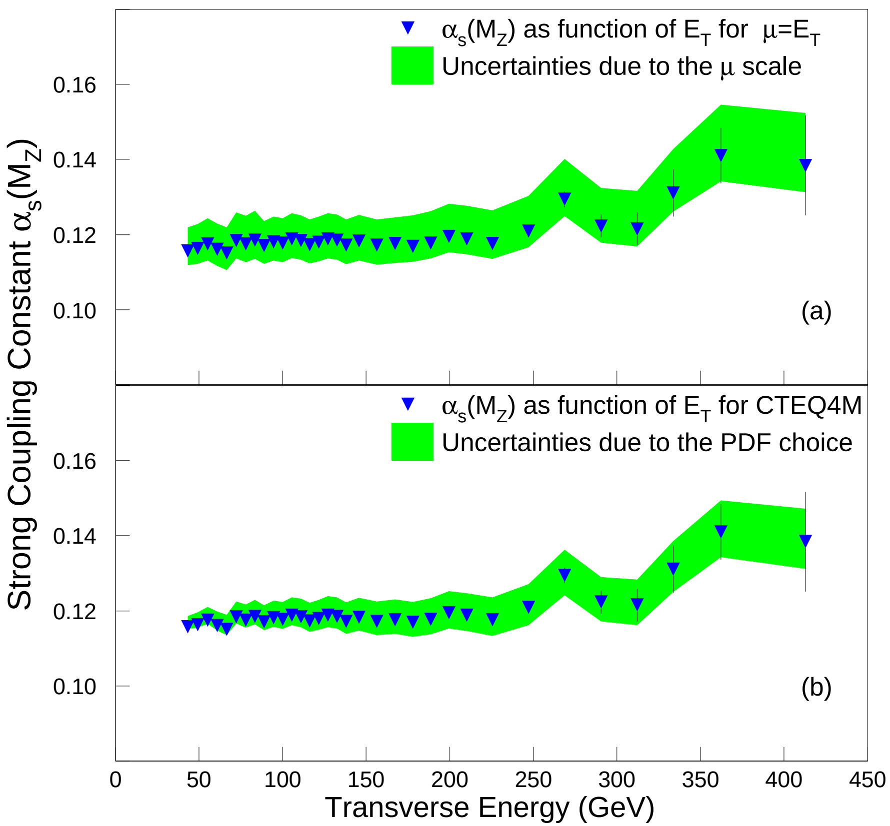
- ✓ all vector boson plus five parton processes numerically

Ellis, Giele, Zanderighi arXiv:0810.2762

Why go beyond NLO?

In many cases, the uncertainty from the pdf's and from the choice of renormalisation scale still give NLO uncertainties that are as big or bigger than the experimental errors.





$$\alpha_s(M_Z) = 0.1178 \ {}^{+6\%}_{-4\%}(scale) \ {}^{+5\%}_{-5\%}(pdf)$$

Why go beyond NLO? - continued

When is NNLO needed?

- ✓ When corrections are large - e.g. H production
- ✓ For benchmark measurements where experimental errors are small

What is known so far?

- ✓ Inclusive cross sections for W, Z and H production

  van Neerven, Harlander, Kilgore, Anastasiou, Melnikov, Ravindram, Smith

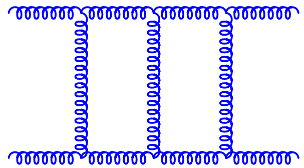
- ✓ Semi-inclusive $2 \to 1$ distributions - W, Z and H rapidity distributions

  Anastasiou, Dixon, Melnikov, Petriello

- ✓ Fully differential $pp \to H, W, Z + X$ — Anastasiou, Melnikov, Petriello
- ✓ DGLAP splitting kernels — Moch, Vermaseren, Vogt
- ✓ NNLO parton distributions — Martin, Stirling, Thorne, Watt

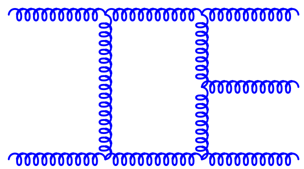
Anatomy of a NNLO calculation

✓ two-loop $2 \to 2$ matrix element



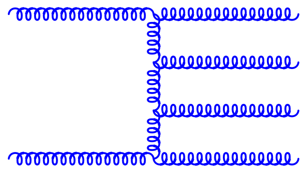
looks like 2 jets in final state

✓ one-loop $2 \to 3$



looks like 2 or 3 jets in final state

✓ tree-level $2 \to 4$

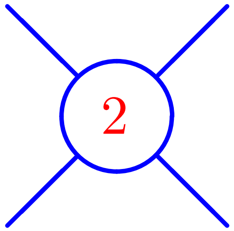


looks like 2, 3 or 4 jets in final state

✓ plus method for combining the infrared divergent parts

Structure of two-loop contribution

✓ The many (thousands) of tensor integrals appearing in two-loop graphs can be written in terms of a few Master Integrals MI_j

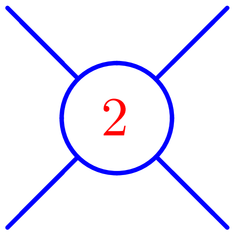


$$= \sum_j a_j \; \mathrm{MI}_j$$

where the a_j are polynomials in kinematical variables and the space-time dimension D.

✓ The MI_j can be expanded in $\epsilon = (4 - D)/2$ so that

$$\sum_{\mathrm{diagrams}}$$



$$= \sum_{i=1}^{4} \frac{X_i}{\epsilon^i} + X_0$$

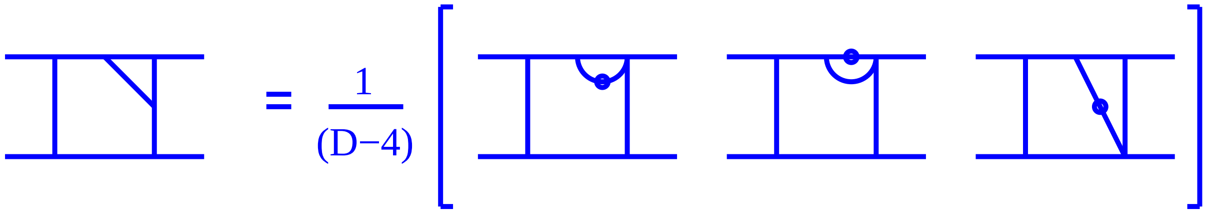
Reduction of Tensor Integrals

✓ Integration by Parts

$$\int d^D k_1 \int d^D k_2 \frac{\partial}{\partial k_i^\mu} \left[\frac{v^\mu}{A_1^{\nu_1} \cdots A_n^{\nu_n}} \right] \equiv 0$$

where v is any momentum in the problem, k_i, p_i.

Chetrykin, Kataev, Tkachov



✓ Packages for solving the IbP identities

- ✓ AIR — Anastasiou, Lazopoulos, hep-ph/0404258
- ✓ IDSolver — Czakon
- ✓ FIRE — A. Smirnov, arXiv:0807.3243

Methods for calculating master integrals

- ✓ Mellin-Barnes contour integrals

  Davydychev; Smirnov; Tausk

- ✓ Differential equations in external scales and match to boundary conditions with fewer scales

  Remiddi, Gehrmann

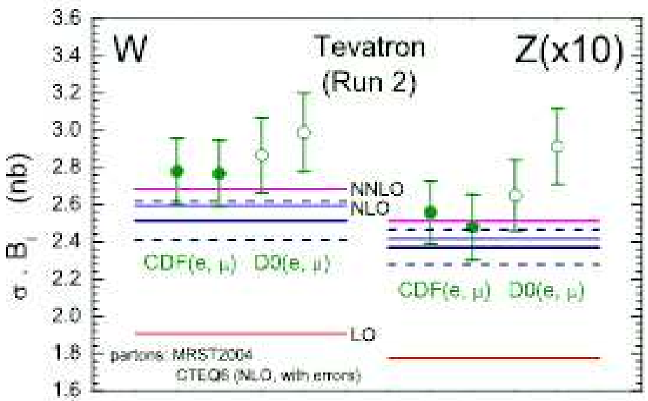
- ✓ Nested sums from Schwinger parameterisation together with Hopf algebra techniques to relate to standard sums

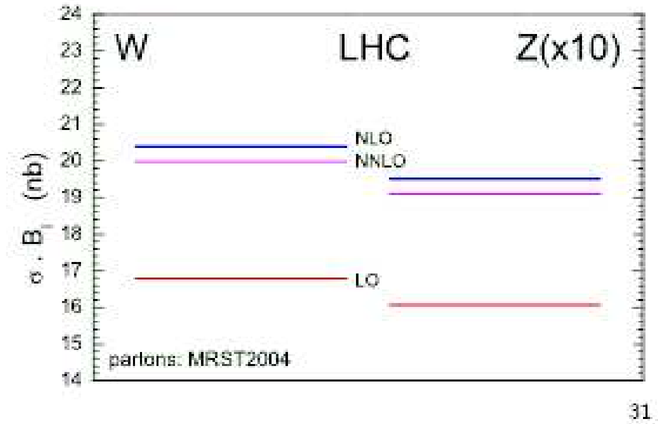
  Moch, Uwer, Weinzierl

- ✓ Numerical method based on iterated sector decomposition - used to check many of above results

  Binoth, Heinrich

Drell Yan production





Most accurate prediction yet

✓ NNLO splitting functions

✓ NNLO PDF fits

✓ NNLO Drell-Yan cross section

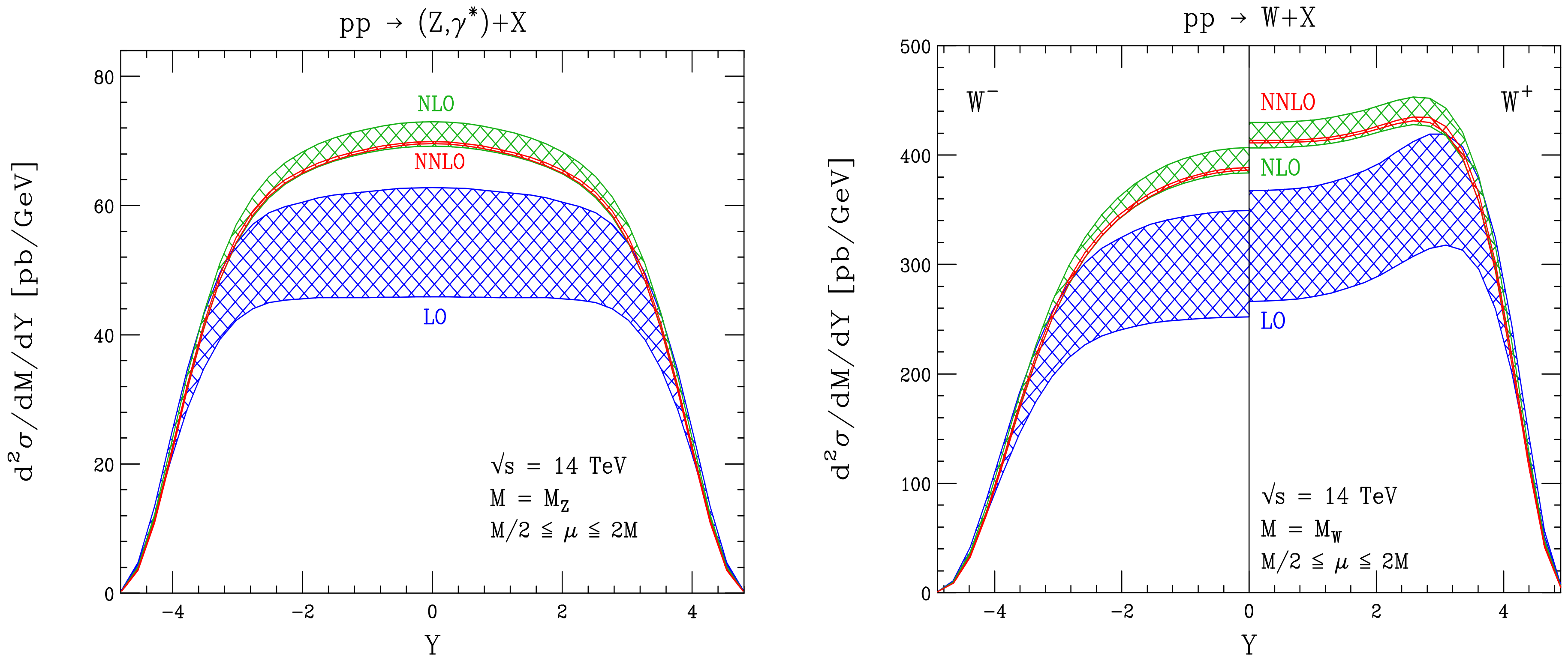
⟹ High precision

Total error of $4\% - -5.5\%$

Martin et al

Aim to able to use as Standard Candle for luminosity measurements.

Gauge boson production at the LHC



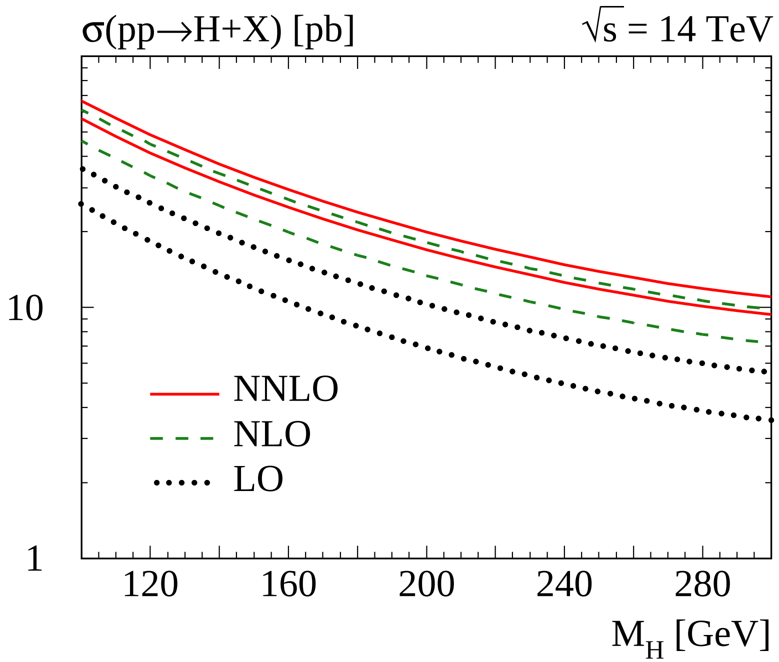
Gold-plated process

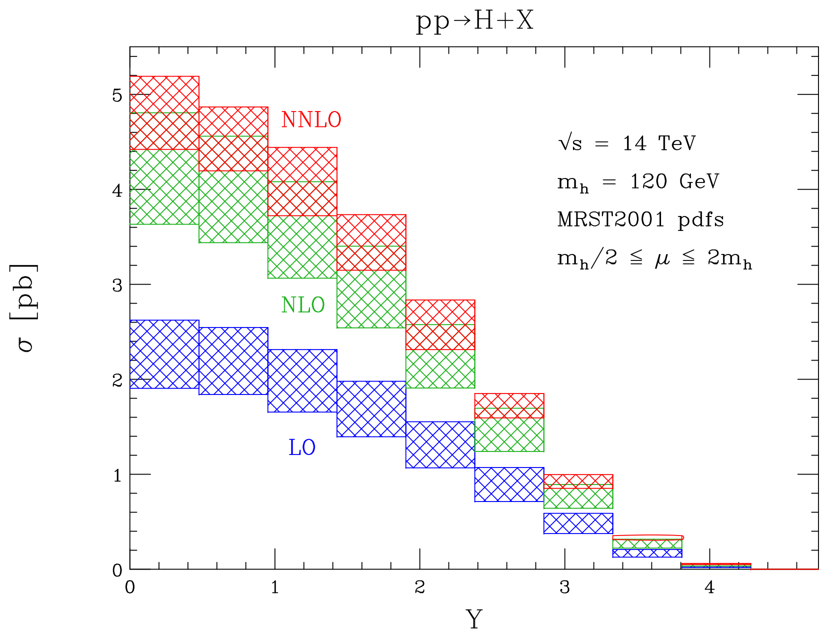
Anastasiou, Dixon, Melnikov, Petriello

At LHC NNLO perturbative accuracy better than 1%

$\Rightarrow$ use to determine parton-parton luminosities at the LHC

Higgs boson production at the LHC





Total cross section

Harlander, Kilgore; Anastasiou, Melnikov, Petriello; . . .

Fully differential

Anastasiou, Melnikov, Petriello

NNLO needed for reliable predictions

Higgs boson production at the LHC

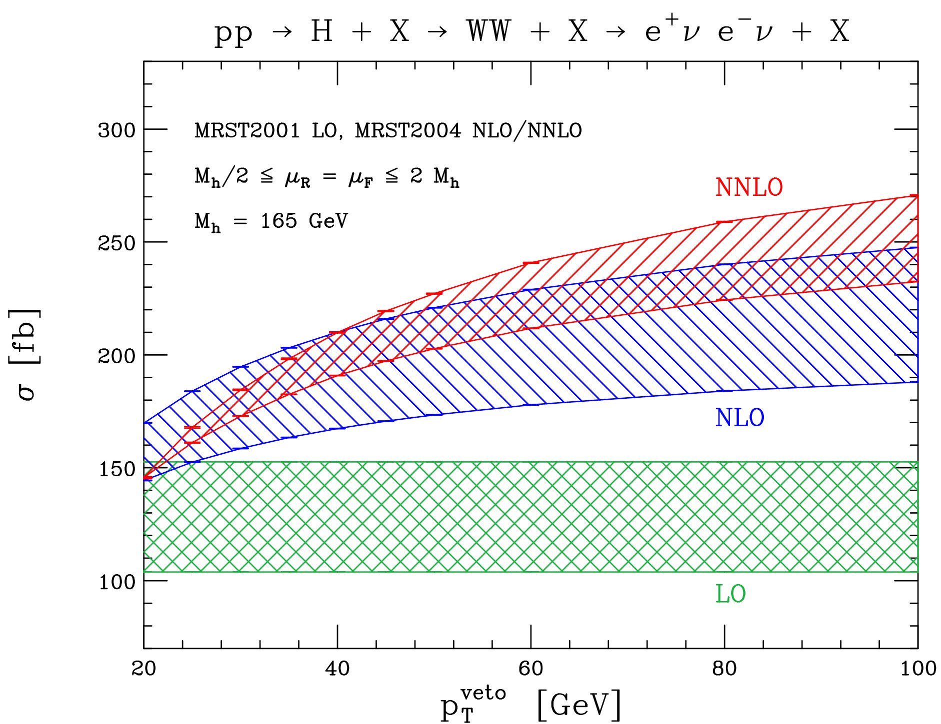
✓ First study of fully inclusive $pp \to H \to WW \to \ell\nu\ell\nu$ with $m_H \sim 165$ GeV

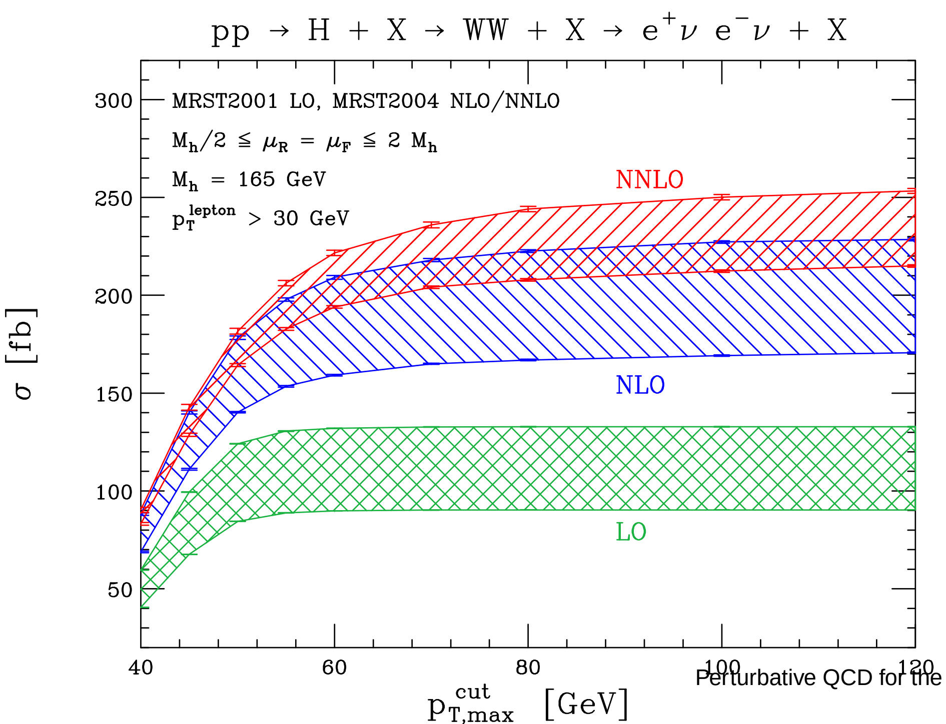
Anastasiou, Dissertori, Stöckli, arXiv/0707.2373

✓ Apply experimental cuts to reduce backgrounds from $t\bar{t}$, non-resonant W^+W^- production

✓ Cuts affect LO/NLO/NNLO cross sections differently

$\Longrightarrow$ shouldn't use inclusive K-factor





Other NNLO calculations on horizon

- ✓ $pp \to jet$ +X
  - ✓ needed to constraint PDF's and fix strong coupling
  - ✓ matrix elements known for some time

    Anastasiou et al, Bern et al
  - ✓ antenna subtraction terms worked out

    Daleo, Gehrmann, Maitre
- ✓ $pp \to t\bar{t}$
  - ✓ necessary for precise m_t determination
  - ✓ matrix elements partially known

    Czakon, Mitov, Moch; Bonciani, Ferroglia, Gehrmann, Studerus, Maitre
- ✓ $pp \to VV$
  - ✓ signal: to study the gauge structure of the Standard Model
  - ✓ background: for Higgs boson production and decay in the intermediate mass range
  - ✓ large NLO corrections

    Chachamis, Czakon, Eiras

Summary

QCD A lot still to do, but progress being made towards main targets

LO largely solved (plus BSM models)

✓ high multiplicity merged with parton shower, ALPGEN, SHERPA, . . .

✗ large theoretical uncertainty

NLO QCD corrections generally large 30% – 100% - much larger than scale variation suggests

✓ Cuts tend to spoil use of inclusive K-factor

✓ Serious effort on Les Houches NLO wish list, several new NLO calculations WWj, $t\bar{t}j$, VVV, $t\bar{t}Z$

✗ ✓ $2 \rightarrow 4$ barrier starting to be breached $q\bar{q} \rightarrow t\bar{t}b\bar{b}$

✓ Numerical methods becoming powerful

NNLO Inclusive and exclusive results for H, W, Z production

✓ DGLAP splitting kernels $\Longrightarrow$ NNLO PDF fits

✗ ✓ $2 \rightarrow 2$ calculations coming onto horizon